

## VERTICAL STRUCTURE OF A SUPERNOVA-DRIVEN TURBULENT, MAGNETIZED INTERSTELLAR MEDIUM

ALEX S. HILL<sup>1,7</sup>, M. RYAN JOUNG<sup>2,3</sup>, MORDECAI-MARK MAC LOW<sup>3</sup>, ROBERT A. BENJAMIN<sup>4</sup>, L. MATTHEW HAFFNER<sup>1</sup>,  
CHRISTIAN KLINGENBERG<sup>5</sup>, AND KNUT WAAGAN<sup>6</sup>

<sup>1</sup> Department of Astronomy, University of Wisconsin-Madison, Madison, WI, USA; alex.hill@csiro.au

<sup>2</sup> Department of Astronomy, Columbia University, New York, NY, USA

<sup>3</sup> Department of Astrophysics, American Museum of Natural History, New York, NY, USA

<sup>4</sup> Department of Physics, University of Wisconsin-Whitewater, Whitewater, WI, USA

<sup>5</sup> Department of Mathematics, Würzburg University, Emil Fischer Strasse 30, Würzburg, Germany

<sup>6</sup> Department of Applied Mathematics, University of Washington, Seattle, WA, USA

Received 2011 September 2; accepted 2012 January 31; published 2012 April 20

### ABSTRACT

Stellar feedback drives the circulation of matter from the disk to the halo of galaxies. We perform three-dimensional magnetohydrodynamic simulations of a vertical column of the interstellar medium with initial conditions typical of the solar circle in which supernovae drive turbulence and determine the vertical stratification of the medium. The simulations were run using a stable, positivity-preserving scheme for ideal MHD implemented in the FLASH code. We find that the majority ( $\approx 90\%$ ) of the mass is contained in thermally stable temperature regimes of cold molecular and atomic gas at  $T < 200$  K or warm atomic and ionized gas at  $5000 \text{ K} < T < 10^{4.2} \text{ K}$ , with strong peaks in probability distribution functions of temperature in both the cold and warm regimes. The  $200\text{--}10^{4.2} \text{ K}$  gas fills  $50\%$ – $60\%$  of the volume near the plane, with hotter gas associated with supernova remnants ( $30\%$ – $40\%$ ) and cold clouds ( $<10\%$ ) embedded within. At  $|z| \sim 1\text{--}2$  kpc, transition-temperature ( $10^5 \text{ K}$ ) gas accounts for most of the mass and volume, while hot gas dominates at  $|z| > 3$  kpc. The magnetic field in our models has no significant impact on the scale heights of gas in each temperature regime; the magnetic tension force is approximately equal to and opposite the magnetic pressure, so the addition of the field does not significantly affect the vertical support of the gas. The addition of a magnetic field does reduce the fraction of gas in the cold ( $<200 \text{ K}$ ) regime with a corresponding increase in the fraction of warm ( $\sim 10^4 \text{ K}$ ) gas. However, our models lack rotational shear and thus have no large-scale dynamo, which reduces the role of the field in the models compared to reality. The supernovae drive oscillations in the vertical distribution of halo gas, with the period of the oscillations ranging from  $\approx 30$  Myr in the  $T < 200 \text{ K}$  gas to  $\sim 100$  Myr in the  $10^6 \text{ K}$  gas, in line with predictions by Walters & Cox.

*Key words:* Galaxy: disk – Galaxy: structure – ISM: kinematics and dynamics – ISM: structure – magnetohydrodynamics (MHD) – turbulence

*Online-only material:* color figures

### 1. INTRODUCTION

Supernovae regulate the structure and dynamics of the multi-phase interstellar medium (ISM; Cox & Smith 1974; McKee & Ostriker 1977). They produce hot ( $\sim 10^6 \text{ K}$ ) gas that occupies a large, though uncertain, fraction of the volume in the disk and halo. Supernovae drive circulation of gas from the disk to heights of several kiloparsecs in a so-called Galactic fountain (Shapiro & Field 1976; Bregman 1980). They also drive turbulence (Norman & Ferrara 1996; Mac Low & Klessen 2004), and a supernova-driven turbulent dynamo can amplify a very small magnetic field to a significant fraction of equipartition (Balsara et al. 2004). Turbulent and magnetic energy densities dominate over the thermal energy in many interstellar environments, particularly in cold atomic and molecular gas (McKee 1990; Heiles & Troland 2005), and provide, along with high-energy particles, a crucial component of the vertical pressure that supports the ISM in hydrostatic equilibrium (Boulares & Cox 1990). Thus, supernovae likely provide the momentum that is responsible for the majority of the pressure support for the ISM.

A complete understanding of the interplay between components of the ISM requires modeling of supernova-driven magne-

tized turbulence in a multi-phase medium. To connect the wide range of temperatures ( $10\text{--}10^7 \text{ K}$ ), densities ( $10^{-4}$  to  $10^3 \text{ cm}^{-3}$ ), and size scales (parsecs to kiloparsecs) found in the diffuse ISM, simulations must have enough spatial resolution near the plane to resolve individual supernova explosions early in their evolution so that unphysical radiative cooling does not remove the energy injected by the supernovae. The simulations also must extend to large enough heights that the mass flux out the top of the simulation box is reasonably small, as well as keep track of heating and cooling of the gas. Variable mesh techniques such as nested grid or adaptive mesh refinement (AMR) are essential to resolve the small-scale structure near the plane while covering a sufficient vertical volume of the Galaxy with available computational resources.

Hydrodynamical and magnetohydrodynamical (MHD) studies of this sort in two (Chiang & Prendergast 1985; Rosen & Bregman 1995; Rosen et al. 1996; Koyama & Ostriker 2009a, 2009b) and three (Korpi et al. 1999; Avillez & Breitschwerdt 2004, 2005a; JounG & Mac Low 2006, hereafter JM06; Gressel et al. 2008a, 2008b; JounG et al. 2009, hereafter JMB09) dimensions (3D) have had considerable success in producing the observed cold, warm neutral and ionized (Wood et al. 2010), and hot gas with reasonable distributions of mass between these components of the ISM near the plane of the Galaxy. However, comparisons to observations suggest three inconsistencies.

<sup>7</sup> Present address: CSIRO Astronomy and Space Science, Marsfield, NSW, Australia.

**Table 1**  
Refinement Zones

z  (kpc)	$\Delta x$ (pc)	
	High Resolution	Medium Resolution
$0. <  z  < 0.05$	1.95	3.91
$0.05 <  z  < 0.3$	3.91	$\leq 7.81$
$0.3 <  z  < 1.$	7.81	15.6
$1. <  z  < 10.$	15.6	15.6
$10. <  z  < 20.$	31.2	31.2

First, the observed distribution of emission measures from the warm ionized medium (WIM) is narrower than the models, suggesting that the real ISM is more compressible than the simulations (Hill et al. 2008; Wood et al. 2010). Second, the simulated WIM disk is thinner than observed (Wood et al. 2010). Third, the observed X-ray temperature of halo gas is cooler than predicted by the models (Henley et al. 2010). Because we expect magnetic fields reduce the compressibility of a fluid and provide vertical support, the addition of magnetic fields may help to address all of these inconsistencies.

With these motivations, we extend the simulations developed by JMB09 to use MHD in this paper. Our primary aim here is to explore the effect of magnetic fields on the vertical structure of the multi-phase ISM. In future work, we will conduct Monte Carlo photoionization simulations (following Wood et al. 2010) of the models presented here and compare the results to  $H\alpha$ , Faraday rotation, and H I observations. In Section 2, we discuss our simulations. In Section 3, we describe the results of the simulations, including the evolution over time (Section 3.1), the phases of the ISM (Section 3.2), the vertical structure (Section 3.3), the magnetic field energy and pressure (Section 3.4), and the thermal and turbulent pressure (Section 3.5). We discuss our results in Section 4, focusing on the phases of the ISM in Sections 4.1–4.3. We discuss the vertical stratification in Section 4.4, with a comment on the substantial relevance of the Walters & Cox (2001) models to the evolution of the halo in our results in Section 4.4.2. We finish with discussion of the generation of the magnetic field (Section 4.5), the effects of numerical resolution (Section 4.6), and some large-scale effects which are not included in our model (Section 4.7). We summarize our main conclusions in Section 5.

## 2. SIMULATIONS

### 2.1. Model

Our model is an extension of the hydrodynamical model developed by JM06 and JMB09. We use FLASH v2.5, a Eulerian astrophysical code with AMR (Fryxell et al. 2000) in 3D. We employ the positivity-preserving HLL3R MHD Riemann solver based scheme developed and implemented in FLASH by Klingenberg et al. (2007), Waagan (2009), Klingenberg & Waagan (2010), and Waagan et al. (2011). The scheme is stable and preserves positivity even for the Mach numbers approaching 100 and very low values of plasma  $\beta$  ( $< 0.1$ ) found in this problem. The code solves the ideal MHD equations, written here in Gaussian cgs units:

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = 0 \quad (1)$$

$$\frac{\partial \rho \mathbf{v}}{\partial t} + \nabla \cdot \left[ \rho \mathbf{v} \mathbf{v} + \left( P_{\text{th}} + \frac{|\mathbf{B}|^2}{8\pi} \right) \mathbf{I} - \frac{\mathbf{B}\mathbf{B}}{4\pi} \right] = \rho g(z) \hat{\mathbf{z}} \quad (2)$$

**Table 2**  
List of Runs

Name	$B_{x,\text{ini}}$ ( $\mu\text{G}$ )	$T_{\text{halo}}$ (1.15 K)	$\Delta x_{\text{min}}$ (pc)	$\Delta x_{\text{max}}$ (pc)	$t_{\text{start}}$ (Myr)	$t_{\text{end}}$ (Myr)
bx0	0.0	$10^6$	3.91	31.2	0	400
bx0hr	0.0	$10^6$	1.95	31.2	300	340
bx10	1.3	$10^6$	3.91	31.2	0	270
bx50	6.5	$10^6$	3.91	31.2	0	400
bx50hr	6.5	$10^6$	1.95	31.2	300	370
bx100	13.	$10^6$	3.91	31.2	0	260

**Note.** The initial temperature at heights  $|z| \gtrsim 1$  kpc is shown as  $T_{\text{halo}}$ .

$$\frac{\partial E}{\partial t} + \nabla \cdot \left[ \left( E + \frac{|\mathbf{B}|^2}{8\pi} + P_{\text{th}} \right) \mathbf{v} - \frac{(\mathbf{B} \cdot \mathbf{v}) \mathbf{B}}{4\pi} \right] = n\Gamma - n^2\Lambda + S \quad (3)$$

$$\frac{\partial \mathbf{B}}{\partial t} - \nabla \times (\mathbf{v} \times \mathbf{B}) = 0. \quad (4)$$

Equations (1)–(4) represent conservation of mass density  $\rho$ , momentum  $\rho \mathbf{v}$ , energy, and magnetic flux. Here  $\mathbf{I}$  is the  $3 \times 3$  identity matrix, and the notation  $\mathbf{B}\mathbf{B}$  denotes the outer product of  $\mathbf{B}$  with itself. The constraint  $\nabla \cdot \mathbf{B} = 0$  applies. The thermal pressure is  $P_{\text{th}}$  (with a  $\gamma = 5/3$  equation of state), the energy density is

$$E = \frac{\rho |\mathbf{v}|^2}{2} + \frac{P_{\text{th}}}{\gamma - 1} + \frac{|\mathbf{B}|^2}{8\pi}, \quad (5)$$

and the magnetic field is  $\mathbf{B}$ . The diffuse heating rate  $n\Gamma(z, T)$ , the cooling rate  $n^2\Lambda(T)$ , and the impulsive heating due to supernovae and stellar winds in superbubbles  $S(\mathbf{x}, t)$  are as specified in JM06 and summarized below. We discuss the gravitational acceleration  $g(z)$  below;  $\hat{\mathbf{z}}$  is a unit vector. Although we do not include explicit Ohmic resistivity in the model, numerical resistivity is present.

The simulation consists of an elongated box  $1 \times 1 \times 40$  kpc<sup>3</sup> in size. The midplane is at the center, with the box extending to  $z = \pm 20$  kpc. We employ zero-gradient boundary conditions (as described in Section 2.1.2 below) at the top and bottom of the box and periodic boundary conditions through the sides. The initial conditions consist of gas in approximate hydrostatic equilibrium. The temperature at the midplane is  $1.15 \times 10^4$  K, while the initial temperature above  $|z| \approx 1$  kpc is  $1.15 \times 10^6$  K. In this paper, we report hydrogen number densities  $n_{\text{H}} \equiv n \equiv \rho/\bar{m}$  adopting a mean mass per hydrogen atom  $\bar{m} = 2.36 \times 10^{-24}$  g, appropriate for a solar metallicity gas. The adaptive mesh is finest near the midplane and becomes coarser with height, as shown in Table 1.

Parameters describing our runs are listed in Table 2. Our naming convention consists of the initial, midplane, horizontal ( $\hat{x}$ ) component of the magnetic field in code units, with hr added for runs with a resolution of  $\approx 2$  pc; runs without the hr label are medium ( $\approx 4$  pc) resolution. We run the medium-resolution setup for 400 Myr. For two models (bx0hr and bx50hr), we also switch to high resolution after 300 Myr. The medium-resolution run establishes an improved initial condition over the hydrostatic equilibrium setup for the high-resolution run at relatively low computational cost. For the majority of this paper, we focus on the results from the bx0hr, bx50hr, and bx100 models, which have initial midplane magnetic fields of 0, 6.5, and 13  $\mu\text{G}$ , respectively. The bx0hr and bx50hr models

ran from 300 Myr to at least 340 Myr, while the bx100 model ran to 260 Myr. Runs were conducted on the Ranger cluster at the Texas Advanced Computing Center; 4 pc resolution runs utilized  $\sim 80$ –200 service units per Myr of evolution, while 2 pc resolution runs required  $\sim 1500$ –2500 service units per Myr.

Supernova explosions are set off in the simulation box as described by JM06, adding  $10^{51}$  erg to a sphere encompassing  $60 M_{\odot}$  centered on the supernova location and redistributing the mass within the sphere. No gas mass is added. This approach prevents excessive cooling in the time steps immediately after the supernova. We use the Galactic supernova rate from JMB09, with Type Ia and core-collapse supernova rates of 6.58 and  $27.4 \text{ Myr}^{-1} \text{ kpc}^{-2}$ , respectively. Three-fifths of the core-collapse supernovae are clustered spatially and temporally to simulate superbubbles. The number of supernovae in each superbubble is drawn from a distribution  $dN_B \propto n_*^{-2} dn_*$  between upper and lower cutoffs  $n_{\text{sn,max}} = 40$  and  $n_{\text{sn,min}} = 7$ , respectively. During the first 5 Myr of each superbubble, additional energy is injected to simulate stellar winds (see JM06). We assign the Type Ia and core-collapse supernovae exponential vertical distributions with scale heights of 325 and 90 pc, respectively. The locations of supernova explosions are random with these restrictions, not correlated with the gas density. The supernova rate does not change over time; this is equivalent to a constant star formation rate. The random times and locations of supernovae are chosen at initialization ( $t = 0$ ), so the supernovae occur in the same locations for all runs with a given configuration independent of resolution: the supernovae occur at the same times and locations in bx0 and bx0hr models, but the supernovae occur at different times and locations in the bx0 and bx50 runs.

As in JM06 and JMB09, we use radiative cooling appropriate for an optically thin, solar metallicity plasma in collisional ionization equilibrium (Dalgarno & McCray 1972; Sutherland & Dopita 1993; see Figure 1 of JM06). We apply a diffuse heating term representing photoelectric heating of dust grains (Bakes & Tielens 1994), thought to be the dominant heating mechanism in the warm and cold neutral medium (WNM and CNM, respectively). The diffuse heating rate is (Wolfire et al. 1995)

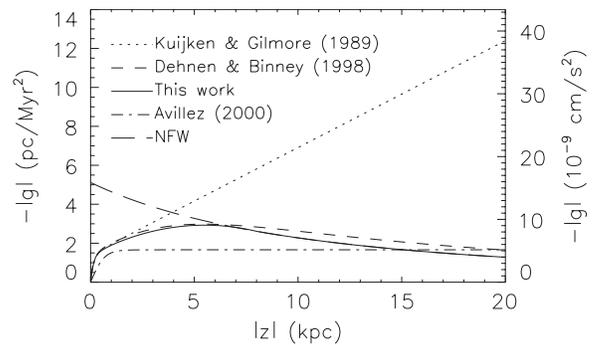
$$\Gamma(z, T) \equiv \begin{cases} \epsilon G_0 e^{-|z|/h_{\text{pe}}} \times 10^{-24} \text{ erg s}^{-1}, & T < 2 \times 10^4 \text{ K} \\ 0, & T > 2 \times 10^4 \text{ K}. \end{cases} \quad (6)$$

We employ a heating efficiency  $\epsilon = 0.05$ , a midplane far-UV (FUV) intensity  $G_0 = 1.7$  times the interstellar value near the Sun (Habing 1968), and an FUV radiation field scale height  $h_{\text{pe}} = 300$  pc.

We do not include any parameterized photoionization heating, as we discuss in Section 2.2.

### 2.1.1. Magnetic Field

The primary modification we make to the physics included in the JM06 and JMB09 models is the addition of magnetic fields. Waagan et al. (2011) present numerous astrophysical tests of the HLL3R MHD solver used in this work. We find that a Courant step of 0.2 is required early in the simulations, as shocks are first established, to obtain a stable solution. After  $\sim 10$  Myr, we increase the Courant step to 0.4–0.8 to improve performance and find that the resulting solutions are generally stable. Densities and temperatures always remain positive by construction; unstable solutions are typically alleviated by manually reducing the time step.



**Figure 1.** Gravitational acceleration using potentials from Kuijken & Gilmore (1989, as used by JM06 and JMB09), Dehnen & Binney (1998) with a spheroidal dark matter halo, Avillez (2000, assuming that the  $\rho_*$  in their Equation (2) should be  $\rho_{*,0}$ ), and Navarro et al. (1996, NFW profile) (normalized to match Equation (7) at  $|z| = 8$  kpc). The potential used in this work (described in Section 2.1.3) is shown with a solid line.

The magnetic field is established with an initially uniform, horizontal field with a constant plasma  $\beta$ , i.e.,  $B_{x,\text{ini}}(z) = B_{x,\text{ini}}(z=0) \times [n(z)/n(z=0)]^{1/2} \equiv B_{x,\text{ini}}$ . Although the random component of the magnetic field in the solar neighborhood is thought to be  $\sim 3$  times the strength of the uniform field (e.g., Ferrière 2001), we do not initialize any random component, allowing it to be generated by turbulence. The initial, midplane magnetic field strengths for each of our runs are listed in Table 2.

### 2.1.2. Box Size

We find that the vertical extent of the box is the most important numerical factor in establishing a realistic temperature distribution in the halo. In particular, the Galactic fountain flow is not accurately captured in simulations presented by JMB09 that only extend to  $z = \pm 5$  kpc (cf. Avillez & Breitschwerdt 2004) and do not enforce outflow at the boundary. In these runs, hot gas from Type Ia supernovae above  $|z| \sim 500$  pc in the first few Myr of the simulation moves rapidly upward. Pressure and density from the top row of the box ( $z = \pm 5$  kpc) are copied to the ghost cells just outside the box. When the pressure inside later drops, after passage of the blast wave, a constant inflow of hot, high-pressure gas is established, unphysically creating ram pressure that heats the gas at  $|z| \gtrsim 1$  kpc to temperatures of  $\sim 10^7$  K. Joung et al. (2012) discuss this numerical issue further. We extend our box to heights  $|z| = 20$  kpc with very low (32 pc) resolution at  $|z| > 10$  kpc. With this box size, we track enough of the mass and energy flux in the fountain flow so that little mass actually flows off the grid.

We employ a gas surface mass density of  $13.2 M_{\odot} \text{ pc}^{-2}$ , the approximate value at the solar circle (Flynn et al. 2006). The star formation rate inferred from the supernova rate we use does not lie on the Kennicutt–Schmidt law (Kennicutt 1989; see also JMB09).

### 2.1.3. Gravitational Potential

We use a modified version of the Kuijken & Gilmore (1989) gravitational potential, consisting of contributions from a stellar disk and a spherical dark halo:

$$g(z) = -\frac{a_1 z}{\sqrt{z^2 + z_0^2}} - a_2 z + a_3 z |z|, \quad |z| \leq 8 \text{ kpc}. \quad (7)$$

The constants  $a_1 = 1.42 \times 10^{-3} \text{ kpc Myr}^{-2}$ ,  $a_2 = 5.49 \times 10^{-4} \text{ Myr}^{-2}$ , and  $z_0 = 0.18$  kpc. In Figure 1, we show the unmodified ( $a_3 = 0$ ) Kuijken & Gilmore (1989) gravitational

acceleration used by JM06 and JMB09. We also show the acceleration due to a parameterized model of the mass distribution of the Milky Way, fitted to the available data by Dehnen & Binney (1998), at the solar circle (Galactocentric radius  $R = R_0 = 8$  kpc). Because of the limited sample of observed stellar radial velocities, the shape of the dark matter halo is uncertain. However, the linear acceleration of the Kuijken & Gilmore (1989) potential (dotted line in Figure 1) overestimates the observed (Dehnen & Binney 1998) acceleration for a flat halo (short dashed line) at  $|z| \gtrsim 4$  kpc.

The shape of the dark halo of the Galaxy is unknown; claims of an oblate, spheroidal, and prolate halo have all been made (e.g., Helmi 2004; Johnston et al. 2005; Fellhauer et al. 2006; Law & Majewski 2010). Therefore, we add the quadratic term in Equation (7) to the acceleration derived from Kuijken & Gilmore (1989) by setting  $a_3 = 5 \times 10^{-5} \text{ kpc}^{-1} \text{ Myr}^{-2}$ , yielding a value of  $g(z)$  roughly consistent with the Dehnen & Binney (1998) constraints in the range  $|z| < 8$  kpc. Note that this functional form is unphysical at heights  $|z| \gtrsim 8$  kpc. At these heights, the gravitational potential is dominated by the dark matter halo. Above  $|z| = 8$  kpc, we use the acceleration of a Navarro et al. (1996) spherical dark matter profile with a scale length of 20 kpc normalized to match Equation (7) at  $|z| = 8$  kpc. Because of the large uncertainties in the Galactic potential at high  $|z|$ , we do not attempt a more detailed fit.

## 2.2. Missing Components

Our models are designed to incorporate the dominant physical processes at work in the diffuse ISM. As discussed by JM06, we do not include the effects of self-gravity, thermal conduction, differential rotation, or a radial variation in the gravity. We also do not include the effect of high-energy particles (cosmic rays), which provide vertical pressure comparable to the kinetic and magnetic pressure in the solar neighborhood (Boulares & Cox 1990; Ferrière 2001). Stellar feedback is only incorporated through supernovae, the stellar wind during the first 5 Myr of the lifetime of each superbubble, and the diffuse heating term. Smaller-scale, lower energy drivers of turbulence and structure formation such as expanding H II regions (Matzner 2002; Brunt et al. 2009) are not included. We note that Koyama & Ostriker (2009a, 2009b) include H II regions but not supernovae in their global two-dimensional simulations.

We do not track UV radiation, so neither photoelectric heating nor photoionization heating are accounted for properly. The diffuse heating term discussed above approximates photoelectric heating by the FUV radiation field in a time- and space-averaged sense. However, Parravano et al. (2003) showed that the UV radiation field and photoelectric heating are likely highly time variable in the solar neighborhood. We comment on the effects of this in Section 4.1. The absence of UV radiation in our models also means we cannot track photoionization, meaning that we can only distinguish between the WNM and WIM as a post-processing step (Wood et al. 2010). The thermal pressure for a given mass density in the WIM is twice that in the WNM due to the extra particle (electron) associated with each hydrogen atom. The impact of this on the dynamics may be small because the thermal pressure is generally significantly smaller than the turbulent pressure in  $10^4$  K gas, as shown in Section 3.5 below. However, we also cannot include photoionization heating, an important heating source in the WIM (Reynolds et al. 1999). Without photoionization heating, the difference between the thermal pressure in real warm ionized gas and that in our model may be more than the factor of two mentioned above due

to the increased temperature. An increase in the diffuse heating rate (Equation (6)) or a second component with a larger scale height could mimic photoionization heating but is not included. We will discuss the effects of UV radiation on the heating rate further in future work.

## 3. RESULTS

### 3.1. Evolution in the Halo

The evolution of the vertical structure of the bx50 model of the ISM is shown through slices of temperature and density over time in Figure 2. The first supernovae that explode at heights  $|z| \approx 1$  kpc occur in a sufficiently low-density environment that they produce shocks that propagate upward, heating the entire volume. Because of our chosen initial condition, the densities at heights  $|z| \gg 1$  kpc are  $\approx 10^{-6} \text{ cm}^{-3}$ . This is much lower than typical halo densities of  $\sim 10^{-4} \text{ cm}^{-3}$ , which can be estimated from dispersion measures of pulsars in the Magellanic Clouds (e.g., Manchester et al. 2006). Therefore, shocks heat the gas more than would be expected with a more realistic initial condition. The initial shocks escape our box after  $\approx 20$  Myr.

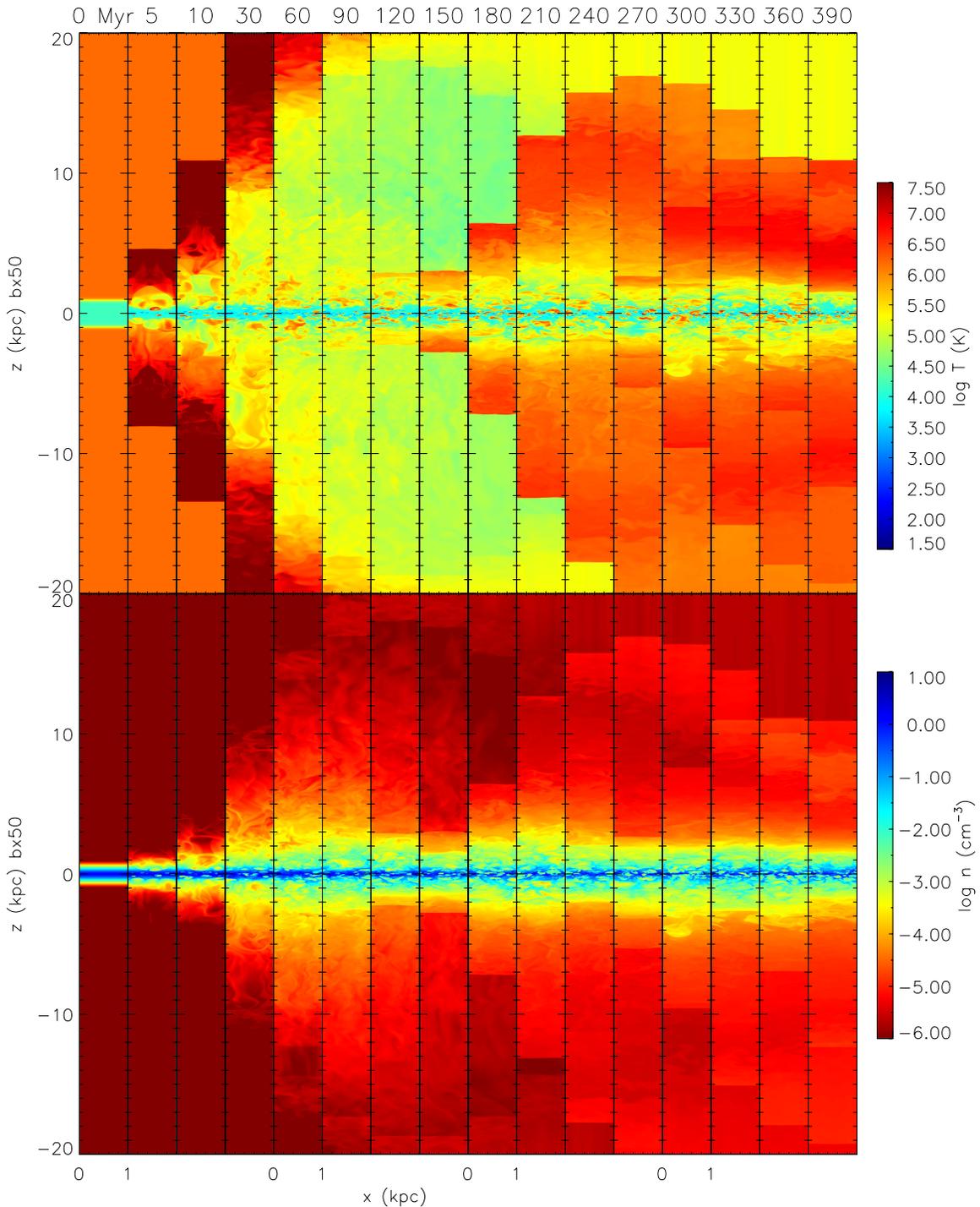
The post-shock gas at very large heights ( $|z| \gtrsim 10$  kpc) cools quickly to temperatures of  $\sim 10^{4.6}$  K, a local minimum in our cooling curve. At lower heights, denser gas is forced upward. This gas is only heated to temperatures of  $\sim 10^5$  K and reaches heights of  $\lesssim 10$  kpc, consistent with a supernova-driven Galactic fountain; we chose our box size to fully capture the cycling of this material.

The fountain material ejected early in our simulations returns to the plane after  $\approx 100$  Myr. At this point, a second pair of shocks propagates away from the plane (see Section 4.4 and Walters & Cox 2001).<sup>8</sup> After this second pair of shocks passes through the medium, the halo density is more consistent with observations, so reasonable initial conditions are now established. Low-resolution simulations confirm that the impact of the initial conditions is negligible after this point. This second pair of shocks heats gas in the halo ( $|z| \gtrsim 3$  kpc) to temperatures of  $\sim 10^6$  K. A layer from 0.5 kpc  $\lesssim |z| \lesssim 2$  kpc consists primarily (by volume) of thermally unstable,  $10^5$  K gas, while warm ( $\sim 10^4$  K) gas dominates (by volume) closer to the plane. We discuss these filling fractions as a function of height below.

### 3.2. Gaseous Components of the ISM

We show vertical and horizontal slices of a snapshot of the bx0hr, bx50hr, and bx100 simulations in Figures 3 and 4. The qualitative structure of each of the models is similar, replicating the conclusion of Avillez & Breitschwerdt (2005a) that the magnetic field does not prevent the breakout of superbubbles in 3D simulations. Mass-weighted histograms of temperature in the plane (black lines in Figure 5) show that the classically defined phases of the ISM at thermally stable temperatures (Wolfire et al. 1995, 2003) do exist in our model: there are peaks in the histograms at  $\approx 25$  K and  $\approx 7000$  K corresponding to cold and warm thermally stable gas. Comparison of mass- and volume-weighted histograms of temperature (Figures 5 and 6) shows that a considerable portion of the volume consists of very low density hot gas: there is a local maximum in the volume-weighted histogram of midplane temperature at  $\approx 10^6$  K (black lines in Figure 6) but no corresponding maximum in the mass-weighted histogram (black lines in Figure 5).

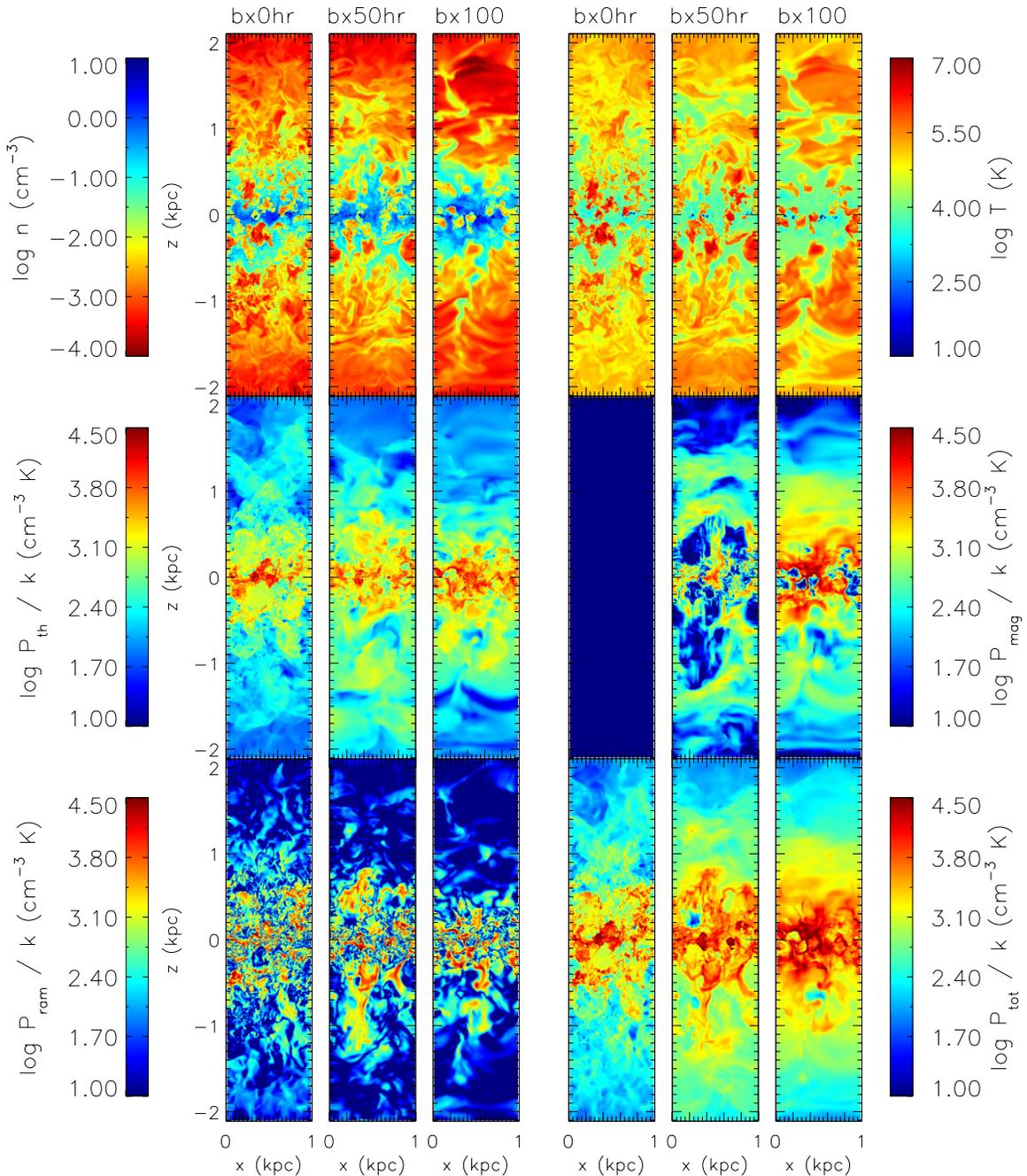
<sup>8</sup> We emphasize that our supernova rate is constant with time, so new star formation is not triggered by the falling fountain material.



**Figure 2.** Slices of temperature (top) and density (bottom) over time for the bx50 model. The time of each slice (in Myr) is listed above each pair of images. Note that the aspect ratio is not unity; the vertical scale of these images is contracted. The horizontal ( $x$ ) axis values are labeled for four slices only to reduce clutter. (A color version of this figure is available in the online journal.)

Figures 5 and 6 also show histograms of temperature and density at  $\approx 500$  pc (red lines) and  $\approx 1$  kpc (green lines) above the plane. The distribution of mass is different at these heights. No cold or cool gas is present. The peak corresponding to warm gas is considerably narrower than in the plane, with the peak at  $\approx 1.5 \times 10^4$  K, twice the most probable temperature of warm gas in the plane. There are also local maxima in the transition-temperature portion of the distribution at  $6 \times 10^4$  K, which is a local minimum on our cooling curve (Figure 1 of JM06), and

$2 \times 10^5$  K, which is thermally unstable and near the maximum cooling efficiency due to O VI. Gas in this transition-temperature regime is normally assumed to lie on interfaces between hot and warm clouds (Fox et al. 2005; Savage & Wakker 2009) because it is highly thermally unstable. However, transition-temperature gas dominates the volume at  $|z| \sim 1$  kpc of our model: the transition-temperature gas occupies  $>90\%$  of the volume between  $|z| = 1.5$  and 2 kpc in the bx50hr model, or  $>95\%$  in the bx0hr model. The transition-temperature gas



**Figure 3.** Vertical snapshots of density, temperature, thermal pressure, magnetic pressure, vertical ram pressure, and total (thermal plus magnetic plus vertical ram) pressure at  $t = 340$  Myr (240 Myr for bx100) from models with a range of initial magnetic field strengths (see Table 2). Slices from each model along the field are shown.

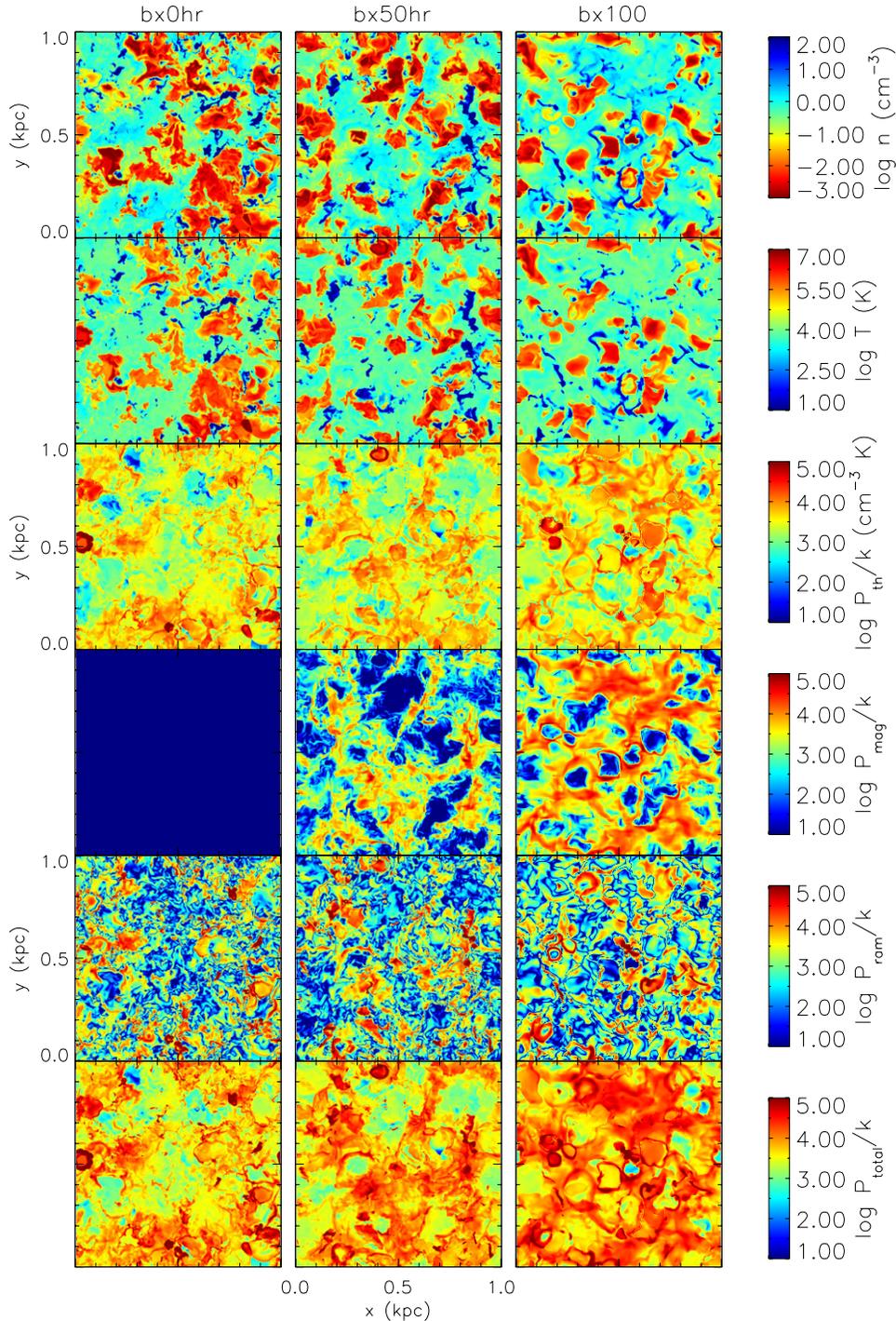
(A color version of this figure is available in the online journal.)

accounts for a similarly large fraction of the mass at these heights as well. The distribution of temperatures within the transition-temperature regime is affected by the magnetic field, with the peak of the distribution (the green lines in Figures 5 and 6) occurring at  $6 \times 10^4$  K in the unmagnetized run but at  $2 \times 10^5$  K in the magnetized runs.

To investigate the role of gas in each temperature regime, we divide the gas into five temperature regimes informed by the three-phase model (McKee & Ostriker 1977): thermally stable “cold” molecular and CNM gas ( $T < 200$  K), thermally unstable “cool” gas ( $200 \text{ K} < T < 5000$  K), thermally stable “warm” gas ( $5000 \text{ K} < T < 10^{4.2}$  K), thermally unstable

“transition-temperature” gas ( $10^{4.2} \text{ K} < T < 10^{5.5}$  K), and “hot” gas ( $T > 10^{5.5}$  K). We use these definitions for the remainder of this paper.

Figure 7 shows the evolution of the surface mass density in the entire box for each temperature regime. After  $\approx 250$  Myr, the gas mass in each temperature regime does not change in the bx0 and bx50 models, indicating that these models have reached a statistical steady state. The bx100 model may have reached a steady state in the distribution of mass as well but has not run for enough time to verify this conclusion. Figure 7 also shows that the formation of cold gas is delayed as the magnetic field strength increases.



**Figure 4.** Images of the midplane slice ( $z = 0$ ) of (top to bottom) density, temperature, thermal pressure, magnetic pressure, vertical ram pressure, and total pressure for a range of initial magnetic field strengths at the same timestamps as in Figure 3. Note that the color scales are different than in Figure 3.

(A color version of this figure is available in the online journal.)

Approximately 62% of the gas mass is in the cold regime (blue lines in Figure 7) in the bx0 model, with 57% of the mass in the cold phase in the bx50 model. The standard deviations of these mass fractions over the time interval 300–400 Myr are  $\approx 0.6\%$ . The mass fractions in the high-resolution (1.95 pc, shown with thick lines) runs are identical to those in the medium-resolution (3.91 pc) ones, suggesting that our results have converged. The steady state mass fractions in warm gas are 28% and 31% for the bx0 and bx50 models, respectively. After 260 Myr, the bx100

model has mass fractions of 48% for the cold and 39% for the warm gas.

We show volume-filling fractions of each temperature regime in Figure 8. In all cases, the warm, thermally stable gas fills  $\approx 36\%$ – $50\%$  of the volume near the plane ( $|z| < 125$  pc), more than any other phase; this fraction is highest in the strongly magnetized case and lowest in the unmagnetized case. Thermally unstable cool gas occupies  $\approx 15\%$  of the volume. Hot and transition-temperature gas ( $T > 10^{4.2}$  K)

**Table 3**  
Experimental Setup

Parameter	Equation	Scale Height (pc)
Gravity ( $z \ll z_0 = 180$ pc)	(7)	8.77
Core-collapse supernovae		90.
Type Ia supernovae		325.
Diffuse heating	(6)	300.

**Notes.** Characteristic heights of physical processes imposed upon the simulation (described in Section 2). All of these values are the same for all runs and do not change over the course of the simulations.

fills most of the remaining 30%–40% of the volume, with cool gas occupying  $\approx 10\%$  in all cases. This hot gas consists primarily of superbubbles produced by correlated supernovae. Note that the evolved picture seen here as well as in similar, recent work (e.g., Avillez & Breitschwerdt 2004; JM06) is one of a pervasive, warm intercloud medium with embedded supernova remnants and superbubbles. This contrasts with the original three-phase ISM models (Cox & Smith 1974; McKee & Ostriker 1977) in which hot gas forms the intercloud medium.

### 3.3. Vertical Structure

We show the space-averaged gas density as a function of height for each temperature regime in Figure 9. Vertical profiles of space-averaged density for each of the three models are shown in Figure 10 and of filling fraction in Figure 11. For comparison, we also show syntheses of observational density profiles in Figures 9(a) and 10. The profiles are the sum of a molecular gas profile with a Gaussian scale height of 81 pc (Clemens et al. 1988; Ferrière 2001), a Dickey & Lockman (1990) profile of atomic hydrogen, and an exponential profile with a scale height of 1.4 kpc representing the WIM (Savage & Wakker 2009).

In our models, cold gas is confined to a narrow layer about the midplane extending to  $|z| \sim 20$  pc. There is no cold gas at heights  $|z| \gtrsim 150$  pc. Even in the plane ( $|z| \lesssim 10$  pc), the volume-filling fraction of the cold gas is very small, 4.2% and 5.2% in the bx0hr and bx50hr models, respectively, as seen in the blue lines in Figure 11. These volume-filling fractions are slightly ( $\sim 10\%$ ) different in the medium-resolution models, as can be seen by comparing the thick and thin blue lines in Figure 8. The images (top rows of Figure 4) show that cold gas primarily exists in filamentary structures in rough thermal pressure equilibrium with the surrounding warm gas.

The scale heights of the warmer phases are larger, with warm gas dominating by volume and mass to  $\sim 500$  pc, above which the transition-temperature gas is the dominant phase. The gas is essentially entirely hot above  $|z| \approx 3$  kpc, with the exception of some transition-temperature gas at  $|z| > 10$  kpc that is above the galactic fountain circulation; this can be seen as the yellow ( $T \approx 10^{5.3}$  K) region at the top of Figure 2 after 210 Myr. Although we include a pressure floor due to the intergalactic medium, we do not model the halo self-consistently. Therefore, we place little physical meaning on the transition-temperature gas at  $|z| > 10$  kpc, above the fountain flow.

Two components of hot gas are evident in Figure 9(f). Near the plane, the hot gas vertical density profile is well described as an exponential distribution that fills  $\approx 20\%$  of the volume. Much of this hot gas has high thermal pressures and low magnetic pressures compared to the surrounding medium (see Section 3.5 below) and is associated with recent supernovae.

**Table 4**  
Scale Heights of Simulated ISM Components

Model	Cold		Cool		Warm		Transition		Hot <sup>a</sup>	
	$h$	$n_0$	$h$	$n_0$	$h$	$n_0$	$h$	$n_0$	$h$	$n_0$
Exponential fits										
bx0hr	8	20	100	0.08	200	0.3	600	0.013	300	0.003
bx50hr	10	9	90	0.09	190	0.3	600	0.012	500	0.001
bx100	11	10	80	0.08	190	0.3	700	0.010	500	0.002
Mass-weighted rms altitudes (Equation (9))										
bx0hr	14		120		260		800		4000	
bx50hr	15		100		260		900		3000	
bx100	17		100		250		800		2000	

**Notes.** Scale heights (in pc) of gas in defined temperature regimes. Fits to an exponential distribution (Equation (8)) are shown in Figure 9. Space-averaged midplane densities  $n_0$  (in  $\text{cm}^{-3}$ ) are listed for exponential fits. Mass-weighted rms altitudes are the average from 320 to 340 Myr.

<sup>a</sup> The exponential distribution of hot gas is fit only to the component near the plane; see the text.

The exponential fit to this distribution finds a scale height of 300–500 pc, larger than the core-collapse supernova scale height of 90 pc. At  $|z| \sim 1$ –3 kpc, the transition-temperature gas dominates, with a second hot component dominating at heights above the transition-temperature gas.

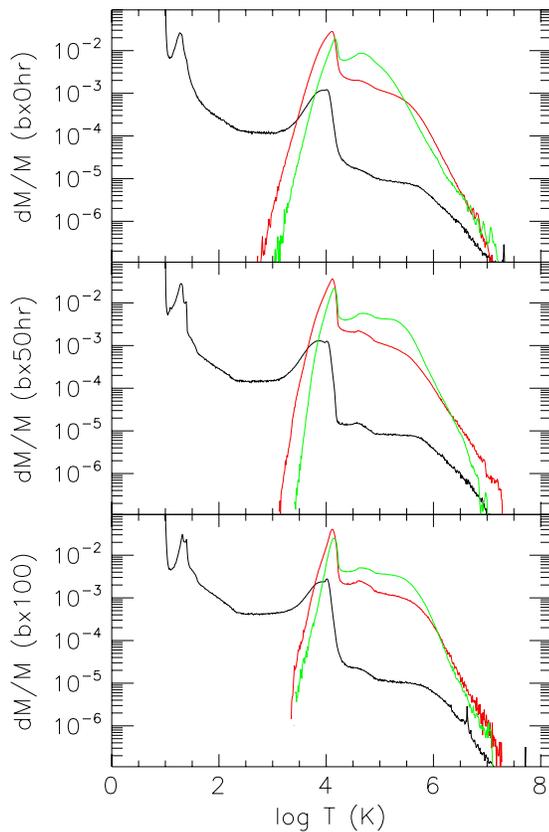
#### 3.3.1. Scale Heights

We characterize the vertical extent of phases of the ISM by calculating characteristic heights of gas in each temperature regime using two methods. The numerical values of the scale heights discussed here are best interpreted by comparison to the physical scales we imposed in our numerical setup, which we describe in Section 2 and summarize in Table 3, rather than directly to observed values.

First, we fit an exponential distribution,

$$n(z) = n_0 \exp\left(-\frac{|z|}{h_{\text{exp}}}\right), \quad (8)$$

to the density profiles shown in Figure 9. Scale heights,  $h_{\text{exp}}$ , and space-averaged midplane densities,  $n_0$ , for each component are listed in Table 4. Each component is fit well by an exponential distribution, though a  $\text{sech}^2(z)$  model may fit the data as well without a cusp at  $z = 0$  (Gómez et al. 2001), and a Gaussian distribution has traditionally been used for the molecular and cold atomic gas (Dickey & Lockman 1990; Ferrière 2001). These space-averaged densities are directly proportional to the total mass in the temperature regime at the height  $|z|$ , although  $n(z)$  is not representative of typical densities within the medium because of the low filling fractions and non-Gaussian distribution of densities (Hill et al. 2008). In particular, the densities of cold gas are much higher than the space-averaged value shown in Figures 9 and 10 due to the very small volume-filling fraction of the cold gas (Figure 11). The exponential distributions of the warm and transition-temperature gas are each truncated at densities  $\sim 10^{-4} \text{ cm}^{-3}$  and heights  $|z| \sim 1.2$ –3 kpc; the fits in Figure 9 and Table 4 were conducted only at heights below the truncation. The halo has a minimum total density due to the imposed pressure floor modeling the presence of the intragroup medium. This floor leads to hot gas at densities of  $\sim 10^{-5.5} \text{ cm}^{-3}$  at  $|z| \gtrsim 5$  kpc. Due to the multiple components of hot gas, we fit its exponential distribution only to hot gas at  $|z| \lesssim 2$  kpc.



**Figure 5.** Mass-weighted histograms of temperature averaged over 21 time slices from  $t = 320\text{--}340$  Myr ( $240\text{--}260$  Myr for bx100) in the bx0hr, bx50hr, and bx100 models (top to bottom). Histograms are calculated within specified height ranges: black lines:  $|z| < 20$  pc; red lines:  $200 \text{ pc} < |z| < 600$  pc; and green lines:  $800 \text{ pc} < |z| < 1200$  pc. Histograms were calculated with logarithmic bin intervals  $d(\log T) = 0.01$ .

Second, we calculate the mass-weighted rms height of gas in each temperature regime (Koyama & Ostriker 2009b):

$$h = \left( \frac{\sum_i (\rho_i z_i^2)}{\sum_i \rho_i} \right)^{1/2}. \quad (9)$$

We performed these estimates with the summations over all zones in each of our temperature regimes; the results are listed in Table 4 and plotted as a function of time in Figure 12. The scale heights fit to an exponential distribution and those derived from Equation (9) are similar for all components except the hot gas. Though it cannot be directly compared to observational determinations of the scale height of the gas, this second technique provides a less biased characterization of the thickness of the medium than the exponential fits because no truncation of the fit is necessary at low densities.

### 3.4. Magnetic Energy

The evolution of the energy contained in the magnetic field is shown in Figure 13. In the bx50 case, the initial magnetic energy density is approximately half the thermal energy density, whereas the initial magnetic energy density in the bx100 case is approximately twice the thermal energy density. In both cases, the thermal energy decreases rapidly and the magnetic energy density increases rapidly in the first 10 Myr. This early behavior is during the initial development of a turbulent cascade

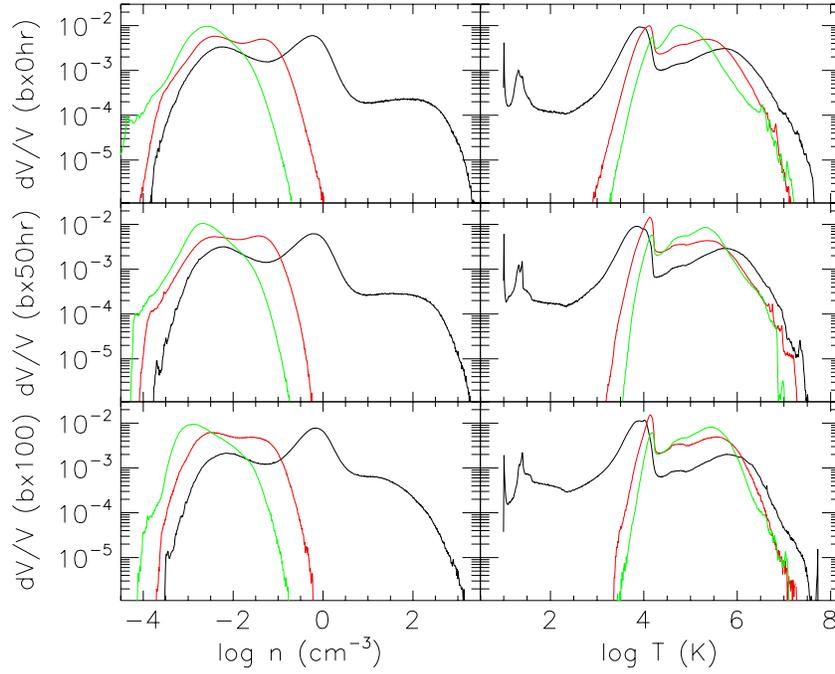
throughout the plane (Balsara et al. 2004). After turbulence is established, the magnetic energy in both the bx50 and bx100 models decreases over time, although there are damped oscillations about the decaying mean with a period of  $\sim 70$  Myr. The field in the bx50 model reaches an approximate steady state at  $\approx 1/3$  of the thermal energy density. We have run the bx100 model for only 260 Myr, and its magnetic field has not reached a steady state. After 260 Myr, the magnetic energy in the bx100 has reduced to approximately the same value as the thermal energy in that model.

### 3.5. Pressure

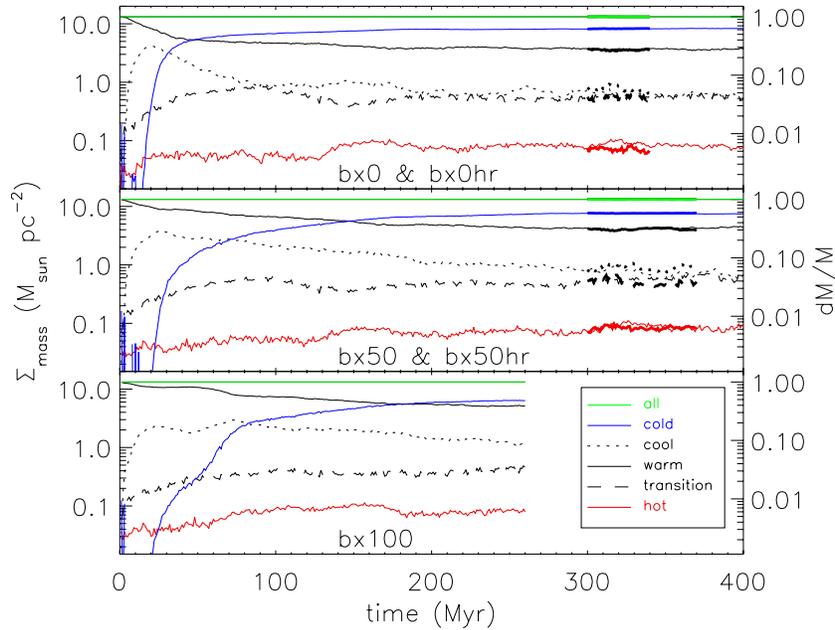
We present images of pressure in Figures 3 and 4. We consider three pressures: the thermal ( $P_{\text{th}} = nkT$ ), magnetic ( $P_{\text{mag}} = |\mathbf{B}|^2/8\pi$ ), and kinetic pressure ( $P_{\text{kin}} = \rho|\mathbf{v}|^2$ ). Plots of these three pressures in the plane as well as their sum are shown as a function of density (phase diagrams) in Figure 14. In the cold (highest density) gas, the non-thermal pressures dominate, with the kinetic pressure being the largest component. The rms velocity in the  $T < 200$  K gas is  $2.5 \text{ km s}^{-1}$ , with typical sonic Mach numbers of  $\sim 5$ . The three pressure components are closest to equipartition in the warm gas ( $n \approx 1 \text{ cm}^{-3}$ ), although the median magnetic pressure in the warm gas is  $\approx 40\%$  of the median thermal or magnetic pressure. The magnetic pressure is very small in the hot gas, where thermal and kinetic pressures are comparable to each other. The total pressure in the cold gas is typically a factor of  $\sim 3\text{--}4$  larger than that found in the warm and hot gas. A small amount of hot gas has thermal pressures  $P_{\text{th}}/k > 10^5 \text{ cm}^{-3} \text{ K}$ , much larger than the total pressures in cold clouds; this overpressured hot gas is typically associated with supernovae which occurred within the last Myr.

In Figure 14(c), we show an estimate of the turbulent pressure, calculated following Equations (2)–(4) of JMB09. In a series of boxes  $8^3 \text{ pc}^3$  in size, we calculate the mass-weighted rms velocity dispersion,  $\sigma_{\text{turb}}$ , in the frame of the center of mass of the box. The turbulent pressure is  $\langle \rho \rangle \sigma_{\text{turb}}^2$ , where  $\langle \rho \rangle$  is the mean mass density within the box. We fit a power-law function to the relationship between density and turbulent velocity dispersion within  $8 \text{ pc}$  boxes,  $\sigma_{\text{turb}} \propto \rho^{-\alpha}$ . We find that a power-law index  $\alpha = -0.35$  best fits the bx50hr model, shallower than the value  $\alpha = -1/2$  reported by JMB09; we find a similar ( $\alpha = 0.35\text{--}0.40$ ) scaling with larger boxes,  $31 \text{ pc}$  on a side, and in the unmagnetized model. Equivalently, instead of the isobaric behavior reported by JMB09, we find that the turbulent pressure increases with density: median turbulent pressures are  $\log(P_{\text{turb},8}/k) = 2.1$  (in  $\text{cm}^{-3} \text{ K}$  units) for the hot gas (defined here by density as  $n < 10^{-1.5} \text{ cm}^{-3}$ ), 2.5 for the warm gas ( $0.1 \text{ cm}^{-3} < \langle n \rangle < 3 \text{ cm}^{-3}$ ), and 3.2 for the cold gas ( $\langle n \rangle > 10 \text{ cm}^{-3}$ ). This trend is in line with that found in the kinetic pressure (Figure 14(d)). For comparison, the turbulent pressure calculated in  $125^3 \text{ pc}^3$  boxes is  $\log(P_{\text{turb},125}/k) = 3.4$ . All  $125^3 \text{ pc}^3$  boxes have mean densities  $\sim 1 \text{ cm}^{-3}$ , so we cannot calculate the turbulent pressure in cold or hot gas on large size scales. Because these estimates capture only the turbulent pressure on size scales of  $\sim 8 \text{ pc}$ , we use the kinetic pressure to calculate the total pressure and for direct comparisons to the thermal and magnetic pressures across widely varying densities.

Considering the thermal pressure by temperature reveals the relative importance of the pressure components in each phase of the ISM. In Figure 15, we show histograms of thermal pressure in the plane for the cold, cool, warm, transition-temperature, and hot gas. The highest thermal pressures are found in hot gas,



**Figure 6.** Volume-weighted histograms of density and temperature for time slices as in Figure 5. As in Figure 5, black lines show the plane, red lines gas at  $|z| \sim 500$  pc, and green lines gas at  $|z| \sim 1$  kpc. Histograms were calculated with logarithmic bin intervals  $d(\log n) = d(\log T) = 0.01$ .



**Figure 7.** Surface mass density as a function of time for the bx0, bx50, and bx100 models (top to bottom). The right axis indicates the corresponding mass fraction. The color and line styles are the same in all plots. Blue lines:  $T < 200$  K. Dotted lines:  $200 \text{ K} < T < 5000$  K. Black solid lines:  $5000 \text{ K} < T < 10^{4.2}$  K. Dashed lines:  $10^{4.2} \text{ K} < T < 10^{5.5}$  K. Red lines:  $T > 10^{5.5}$  K. Green lines: all gas. High-resolution versions (bx0hr and bx50hr) are overplotted (on bx0 and bx50, respectively) with the same line styles and colors but with thick lines.

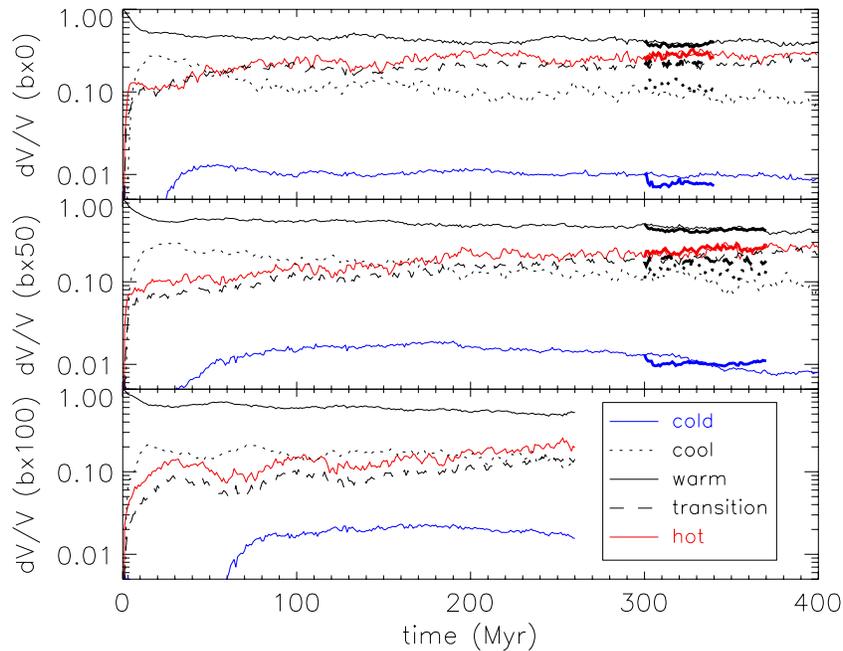
while the lowest are found in cool gas. The range of pressures found in cold gas is smaller, with  $2.5 \lesssim \log(P_{\text{th}}/k) \lesssim 4$ . The distribution of each  $T > 200$  K temperature regime is roughly symmetric about the most probable value of  $\log P_{\text{th}}$ , while the distribution of pressure for the cold gas is highly asymmetric, with a peak at  $P_{\text{th}}/k \approx 10^3 \text{ cm}^{-3} \text{ K}$  but pressures extending to  $10^4 \text{ cm}^{-3} \text{ K}$ . Cool, warm, and transition-temperature gas are all found with thermal pressures both above and below those found in the cold gas. We note that Avillez & Breitschwerdt (2005a, Figure 6) found the cold gas to have a lower thermal pressure than the other phases.

### 3.6. Vertical Support

What provides the vertical support for the stratified ISM in our models? To address this question, we evaluate the roles of vertical kinetic pressure, thermal pressure, magnetic pressure, and magnetic tension (Piontek & Ostriker 2007):

$$P_{\text{tot}}(z) = \langle \rho v_z^2 \rangle + \langle nkT \rangle + \frac{\langle |\mathbf{B}|^2 \rangle}{8\pi} - \frac{\langle B_z^2 \rangle}{4\pi}, \quad (10)$$

where the brackets denote volume averaging over  $x$  and  $y$  positions at fixed  $z$ . The contributions of these vertical pressure



**Figure 8.** Time evolution of volume-filling fractions of phases of the ISM, selected by temperature, in the plane ( $|z| < 125$  pc). Plots are shown for models with three different magnetic field strengths (bx0, bx50, and bx100), with the two high-resolution models (bx0hr and bx50hr) shown with thick lines. Line styles are the same as in Figure 7.

components with height are shown for a snapshot of the bx50hr model in Figures 16(a) and (c) and for a snapshot of the bx0hr model in Figures 17(a) and (c). If the gas is in hydrostatic equilibrium, the pressure at a height  $z$  balances the weight  $W(z)$  of the gas above  $z$ . Because all components of the pressure at the top of our box ( $|z| = 20$  kpc) are negligible, this condition would imply (Piontek & Ostriker 2007)

$$P_{\text{tot}}(z) = \int_z^{z_{\text{max}}} g \rho dz \equiv W(z). \quad (11)$$

We compare the weight and pressure of the ISM in panels (b) and (d) of Figures 16 and 17. As required by the identical surface mass densities and similar scale heights of the magnetized and unmagnetized models (see Section 3.3.1), the total pressure profiles are similar in the two models. In both cases, the kinetic and thermal pressure terms are each significant, with the thermal pressure larger than the kinetic pressure at the midplane. In the magnetized model, the magnetic pressure is smaller than either the kinetic or thermal terms at most heights. The total, large-scale impact of the magnetic field is smaller still because the sum of the magnetic tension and magnetic pressure is very small: the maximum of  $|\mathbf{B}(z)|^2(8\pi)^{-1} - B_z^2(z)(4\pi)^{-1}$  is  $< 10\%$  of the maximum of  $P_{\text{tot}}$ . Equivalently, the mean value of the horizontal field,  $B_x^2 + B_y^2$  is approximately equal to the vertical field,  $B_z^2$ .

#### 4. DISCUSSION

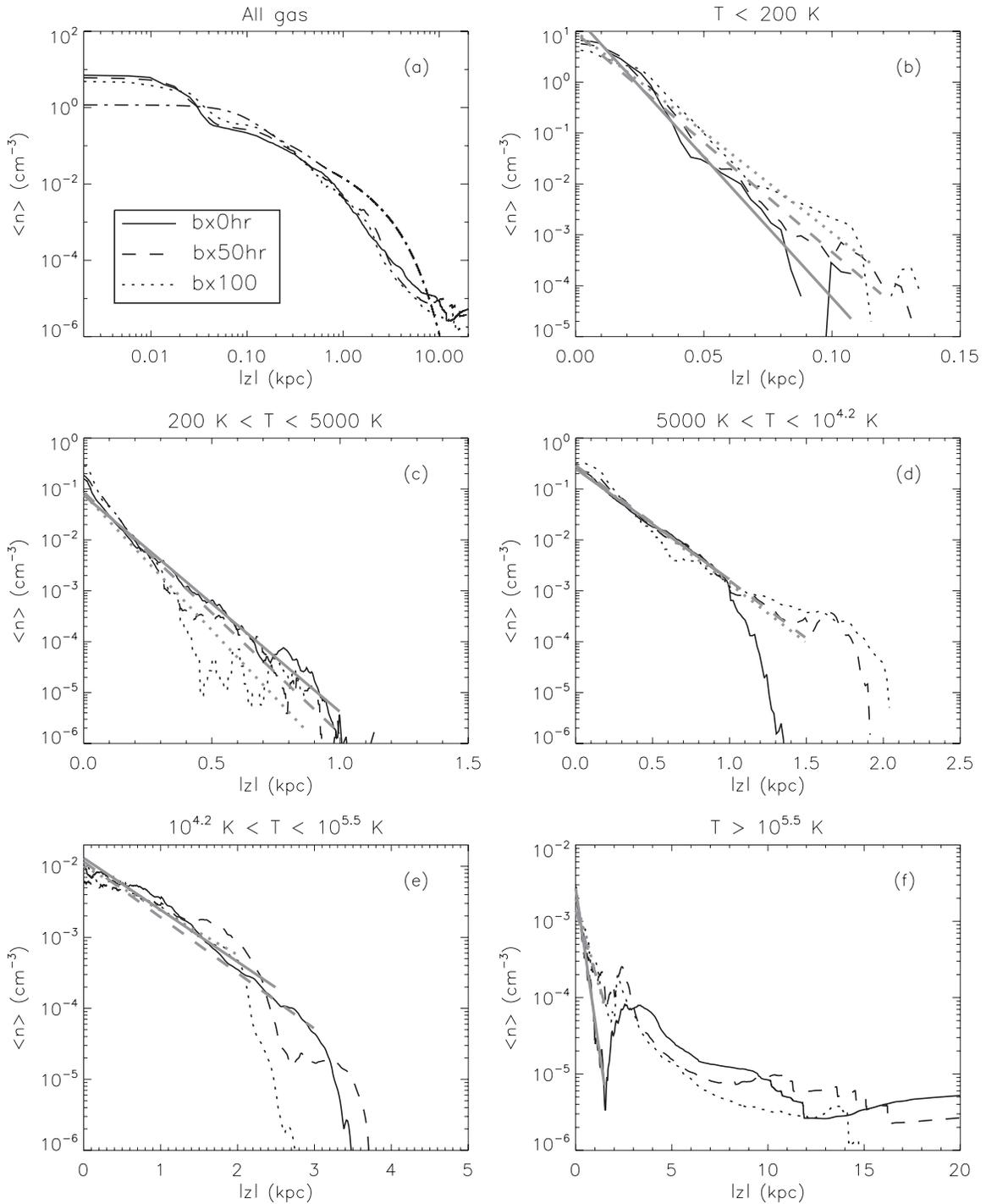
Supernova-driven simulations produce an ISM with mass and height distribution among temperature regimes of the ISM broadly similar to observed values. This suggests that dynamical pressure is the principal driver of vertical stratification of the medium and that supernovae provide the bulk of the dynamical energy.

##### 4.1. Phases and Diffuse Heating

We find strong peaks for warm and cold gas in mass- and volume-weighted histograms of temperature corresponding

to temperature regimes that are thermally stable as determined by the cooling curve (Figures 5 and 6), similar to the magnetorotational instability-driven simulations of Piontek & Ostriker (2007). This contrasts with the results reported by Avillez & Breitschwerdt (2004, 2005a), in which probability density functions (PDFs) of temperature are essentially flat, while Korpi et al. (1999) find weaker but still significant peaks in their simulations (which extend to only  $|z| = 1$  kpc and therefore do not track the Galactic fountain flow). The difference is likely due to the adopted diffuse heating rate. As noted by JM06, we have chosen a diffuse heating rate too small to maintain the hydrostatic equilibrium initial condition absent supernovae because, in reality, the pressure support for hydrostatic equilibrium is maintained by the combination of photoelectric heating (modeled by our diffuse heating term), heating due to supernovae, and non-thermal pressure. Avillez & Breitschwerdt (2004, 2005a), following the example of Mac Low et al. (1989) and Mac Low & Ferrara (1999), used a diffuse heating term that is sufficient to maintain hydrostatic equilibrium by itself. Their diffuse heating term is therefore 18 times larger than that used here. Our diffuse heating rate does allow  $\approx 12\%$  of the total gas mass to cool to the minimum allowed temperature, 10 K, as can be seen in the peak in histograms of temperature (Figures 5 and 6). Of the cells in the simulation that cool to the minimum temperature, 85% (accounting for 95% of the mass at the minimum temperature) have a density  $n > 100 \text{ cm}^{-3}$ .

Proper treatment of the diffuse heating requires the inclusion of radiative transfer to allow a heating rate that is not constant in both time and space. Parravano et al. (2003) argued that the FUV radiation field intensity in the solar neighborhood varies by factors of 2–20 on timescales of  $\sim 100$  Myr. Furthermore, the coldest gas will be heated by ionization from local star formation (e.g., Matzner 2002) in ways that we do not model here. Therefore, a detailed model of the distribution of temperatures into phases must account for these effects.

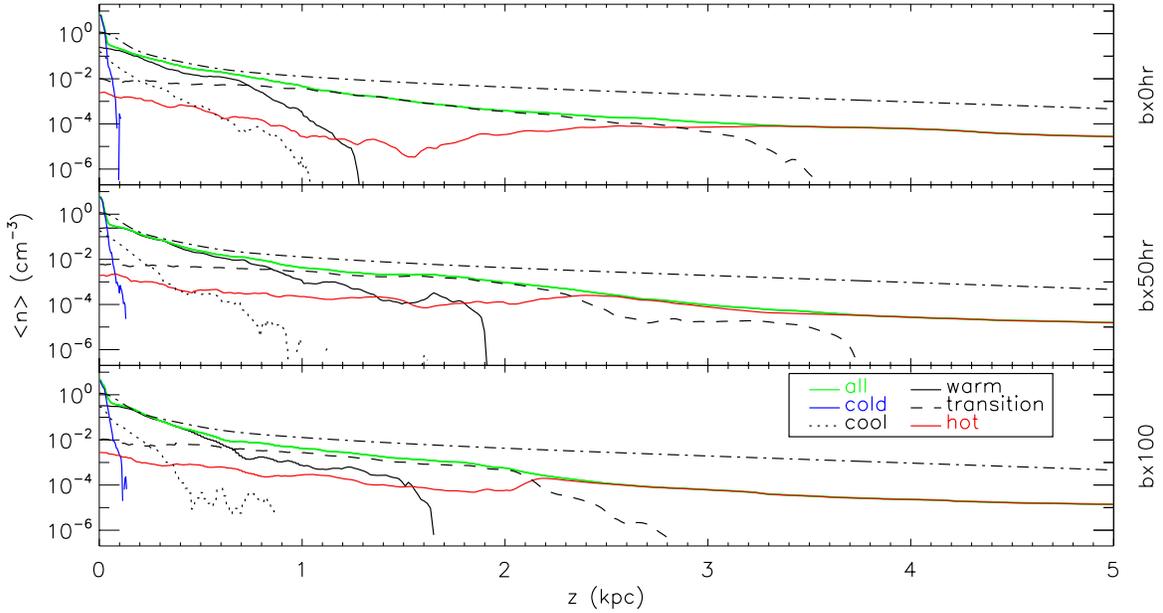


**Figure 9.** Vertical profiles of horizontally averaged density as a function of height. Profiles are shown for all gas on a log–log scale (panel (a)) and on a semi-log scale for the cold, cool, warm, transition-temperature, and hot gas (panels (b)–(f)). Each plot includes dark lines for the `bx0hr`, `bx50hr`, and `bx100` models with data above and below the plane in four, equally spaced time slices from 330 to 360 Myr (230–260 Myr for the `bx100` model) averaged. Exponential fits are shown with gray lines, with scale heights listed in Table 4. Aspect ratios are different in each plot. The dot-dashed line in panel (a) represents the sum of fits of observed  $\text{H}_2$ ,  $\text{H}^0$ , and  $\text{H}^+$  density profiles (see Section 3.3).

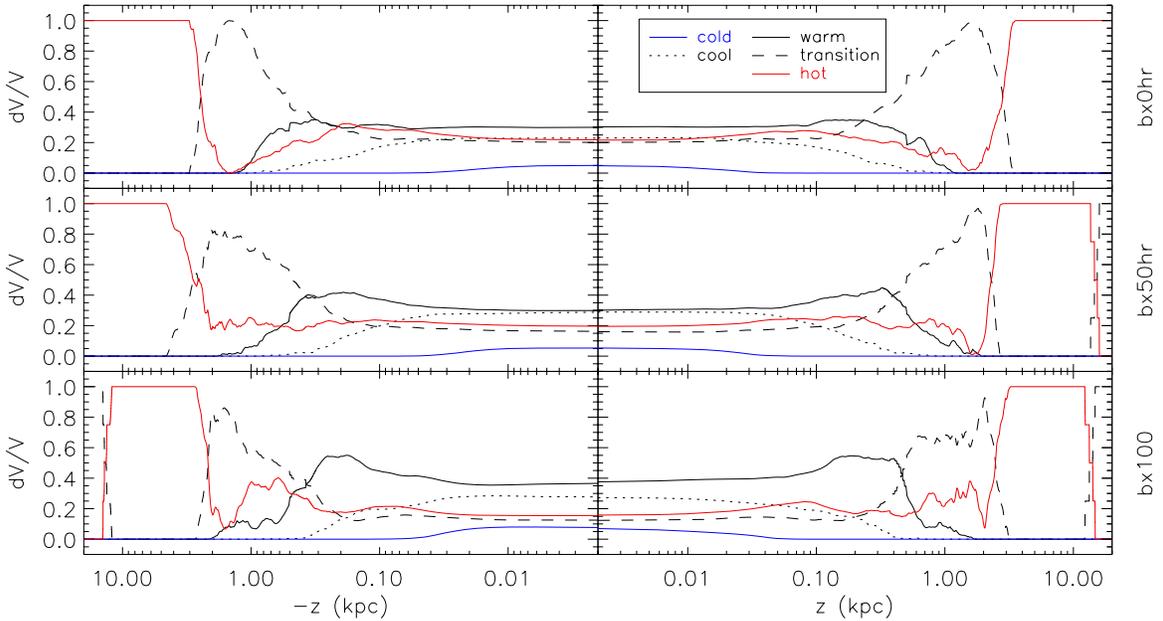
#### 4.2. Filling Fraction of Thermally Unstable Atomic Gas

Comparisons of HI emission and absorption data find that gas in the range  $500 \text{ K} < T < 5000 \text{ K}$  accounts for  $\gtrsim 50\%$  of the WNM column density (Heiles & Troland 2003). We conduct two crude estimates of the fraction of WNM gas in our simulations that is thermally unstable. If we assume that  $\sim 30\%$  of the warm gas in our simulation is ionized (Ferrière 2001) while the cool gas is all neutral, thermally unstable gas accounts

for  $\sim 40\%$  of the mass in the range  $200 \text{ K} < T < 10^{4.2} \text{ K}$  in our magnetized and unmagnetized simulations. This estimate is relatively insensitive to the chosen maximum temperature considered “thermally stable”; we use the cutoff of 5000 K used by Heiles & Troland (2003). Alternatively, 50% of the mass in the  $500 \text{ K} < T < 8000 \text{ K}$  range is cooler than 5000 K; this value is the same in our magnetized and unmagnetized models. We note that Avillez & Breitschwerdt (2005a) found a 60% unstable



**Figure 10.** Horizontally averaged density as a function of height, with the bx0hr, bx50hr, and bx100 models shown separately (top to bottom). The data shown and the dot-dashed line (representing fits to observed molecular, atomic, and ionized hydrogen profiles) are the same as in Figure 9, while the line styles depicting the data are the same as in Figures 7 and 8.



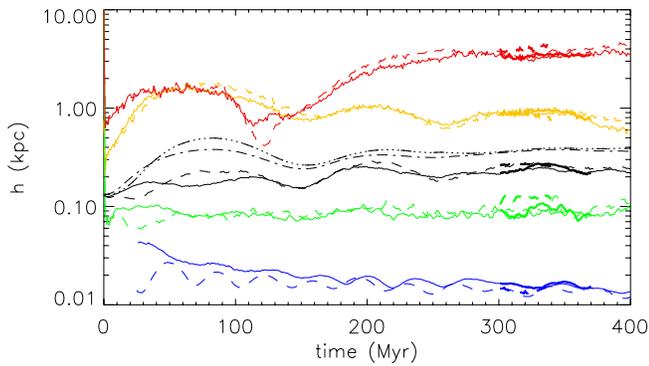
**Figure 11.** Volume-filling fraction as a function of height for the bx0hr, bx50hr, and bx100 models (top to bottom) after 340 Myr (240 Myr for bx100). Line styles and colors are the same as in Figures 7, 8, and 10. Note that the horizontal ( $z$ ) axis scale is logarithmic.

WNM fraction using this estimate, although (as noted above) their gas is roughly evenly distributed across all temperatures whereas ours is concentrated at temperatures  $T \gtrsim 3000$  K.

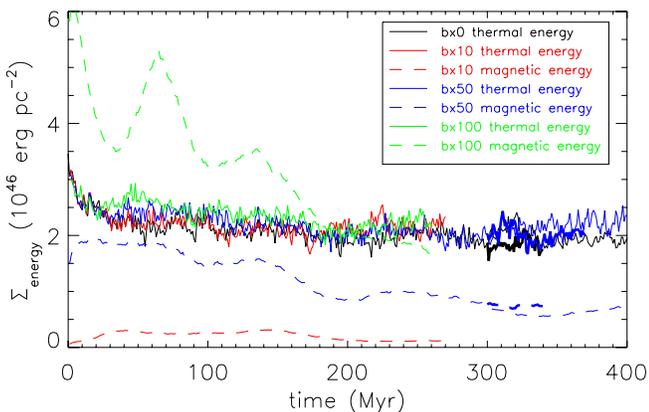
Photoionization modeling promises to allow a cleaner discrimination between WNM and WIM gas. We also plan to construct synthetic H I observations through our simulations to refine this estimate. Although diffuse ionized gas is rarely observed at temperatures  $< 6000$  K (Haffner et al. 1999; Madsen et al. 2006), gas at  $T < 6000$  K in our simulations could be ionized: photoionization heating, the dominant heating source in the WIM at densities  $n \gtrsim 0.1 \text{ cm}^{-3}$  (Reynolds & Cox 1992; Reynolds et al. 1999), would increase the temperature of such gas.

### 4.3. Thermal Pressure of CNM and Molecular Gas

Jenkins & Tripp (2011) have compiled a set of thermal pressures in CNM gas using Space Telescope Imaging Spectrograph C I observations. They find that the CNM pressure is lognormally distributed when weighted by mass. Their estimated volume-weighted distribution, calculated assuming that the CNM gas is in thermal equilibrium (effective  $\gamma = 0.7$ ), is plotted with the dotted blue line in Figure 15. The highest thermal pressure in cold gas in the model is higher than observed in the CNM, while the peak of the modeled distribution is at lower pressure than most of the observed CNM pressures. Although we do not include molecule formation in our model,



**Figure 12.** Scale heights (Equation (9)) of components of the simulated ISM as a function of time. Hot gas is shown with red lines, transition-temperature gas with yellow, warm gas with black, cool gas with green, and cold gas is shown with blue; we listed these models in order of decreasing scale height after 200 Myr. The bx50 and bx50hr models are shown with solid lines, and the bx0 and bx0hr models with dashed lines (with the high-resolution portion of the runs with thick lines in both cases). The cold gas scale height is not plotted before 25 Myr because of noise before significant amounts of cold gas form. The scale height of all gas is shown with a dot-dashed black line for the bx50 model and a triple-dot-dashed black line for bx0.



**Figure 13.** Surface magnetic (dashed lines) and thermal (solid lines) energy density as a function of time. The unmagnetized bx0 and magnetized bx10, bx50, and bx100 models are shown with thin lines, while the high-resolution counterparts (bx0hr and bx50hr) are shown with thick lines.

we now attempt to divide cold atomic gas from molecular gas by setting a density threshold of  $100 \text{ cm}^{-3}$ , above which the timescale for molecule formation is  $\lesssim 10$  Myr (Hollenbach et al. 1971). Of this gas we identify as “molecular,” 51% has cooled to the minimum temperature (10 K). Cold gas at  $n < 100 \text{ cm}^{-3}$ , which we identify as “CNM,” is shown with the dashed blue line in Figure 15. The volume-weighted distribution of cold,  $n < 100 \text{ cm}^{-3}$  gas is similar in width to the Jenkins & Tripp (2011) observations but the most probable value is  $\approx 0.5$  dex lower; our *molecular* ( $n > 100 \text{ cm}^{-3}$ ) gas has a similar median thermal pressure to that observed in the *atomic* gas. However, we note that such dense gas is, in reality, often part of larger, self-gravitating structures that will increase its pressure, physics that is not included in our models.

We also note that, by construction, our “molecular” gas must have a thermal pressure  $P_{\text{th}}/k > 1000 \text{ cm}^{-3} \text{ K}$ , whereas our “CNM” gas at the maximum temperature commonly found of  $\approx 30 \text{ K}$  (see histograms in Figure 5) has a maximum thermal pressure  $P_{\text{th}}/k = 3000 \text{ cm}^{-3} \text{ K}$ . Therefore, we can learn relatively little by dividing molecular from atomic gas by density in our simulations.

With that strong caveat, we investigate the extent to which star-forming clouds can form in our simulations, absent self-

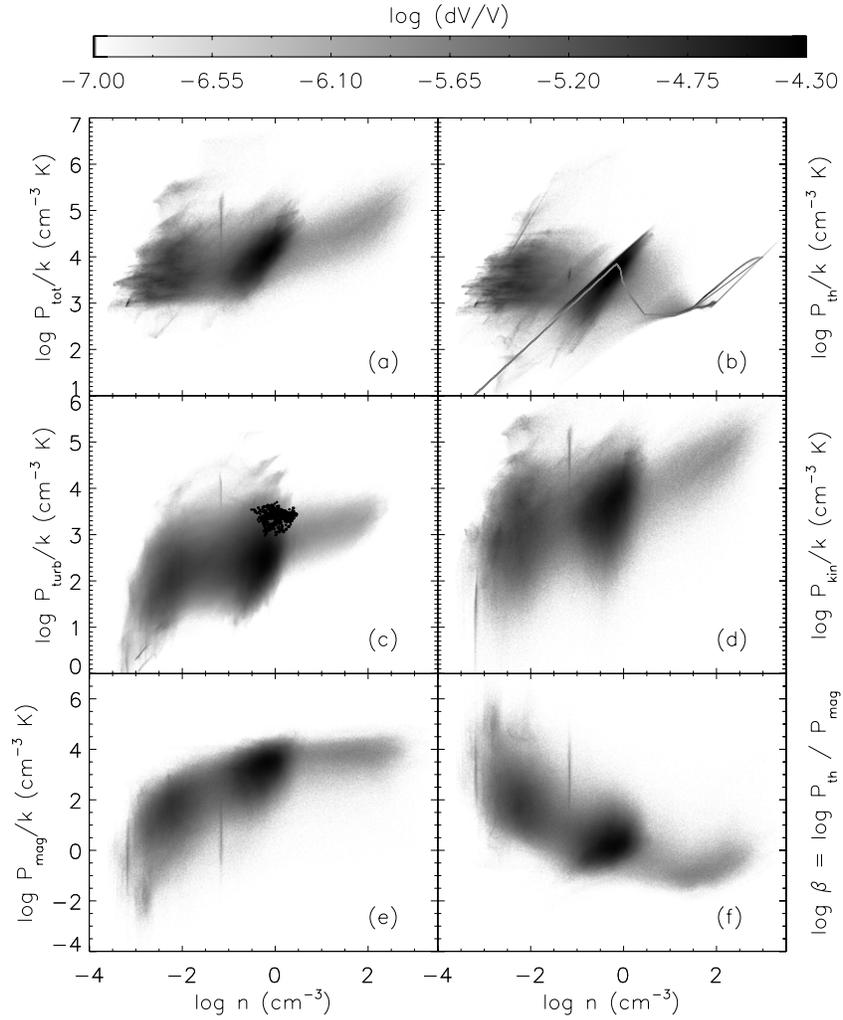
gravity. JM06 argued that the triggered star formation rate in their hydrodynamical simulations was  $\lesssim 10\%$  of the input star formation rate. Kainulainen et al. (2009) found that PDFs of column densities in nearby molecular clouds are generally lognormal, suggesting that turbulent motions dominate. However, active star-forming clouds have power-law wings at high column density. Kainulainen et al. (2011) argue that the non-lognormal tail represents a transition from diffuse clouds dominated by turbulence to gravitationally bound star-forming clouds. Our models show no evidence for non-lognormality, as expected given the lack of self-gravity. The distribution of the logarithm of densities (which we expect to be Gaussian if the column density distribution is lognormal; Vázquez-Semadeni & García 2001) of cold ( $T < 200 \text{ K}$ ) gas in our simulations is symmetric about the median density ( $\langle n \rangle = 43 \text{ cm}^{-3}$ , with a skewness of  $\approx -0.04$ ). This distribution is somewhat flattened compared to a lognormal, with a kurtosis of  $\approx -0.4$  (set such that a Gaussian has zero kurtosis). Our molecular gas is part of a continuous distribution with the cold atomic gas. Therefore, we argue that our simulation does not contain the physics necessary for molecular cloud formation: colliding flows in supernova-driven turbulence alone are insufficient to drive significant star formation.

We now consider the relative roles of thermal and turbulent pressure on small ( $\sim 8 \text{ pc}$ ) and large ( $\sim 125 \text{ pc}$ ) size scales. These pressures are presented in Section 3.5 and shown in Figures 14(b) and (c). On small scales, the turbulent pressure in the warm gas is  $\approx 1$  dex below the thermal pressure in the warm gas, while the large-scale turbulent pressure is comparable to the thermal pressure in the warm gas, with median values of  $\log(P_{\text{th}}/k [\text{cm}^{-3} \text{ K}]^{-1}) = 3.5$  and  $\log(P_{\text{turb},125}/k [\text{cm}^{-3} \text{ K}]^{-1}) = 3.4$ . This result is consistent with the claim of Wolfire et al. (2003) that thermal pressure in the WNM is larger than turbulent pressure on scales smaller than  $\approx 215 \text{ pc}$ . They expect the time between supernova-induced shocks in the plane to be greater than the cooling time, so the CNM and WNM should be in approximate thermal pressure equilibrium even in the presence of supernova-driven turbulence.

#### 4.4. Vertical Stratification

Our models produce a stratified medium. The medium is in approximate hydrostatic equilibrium: the *total* pressure including thermal, magnetic, magnetic tension, and kinetic components of each layer balances the weight of the gas above.

The cold gas is confined close to the plane, with the warmer phases distributed with increasing scale heights. The effects of the magnetic field on vertical stratification are most evident in the cold gas, the only temperature regime in which magnetic pressure dominates over thermal pressure (Section 3.5 and Figure 14). The scale height of the cold gas in the bx50 model (solid blue line in Figure 12) is  $\sim 20\%$  higher than that of the cold gas in the bx0 model (dashed blue line), particularly over the first  $\sim 200$  Myr of the simulation. At late times, the field in the bx50 model has partially dissipated away (see Section 4.5), and the difference between the cold gas scale heights reduces. This can be seen by comparing Figures 12 and 13. In all of the warmer temperature regimes, the scale heights of each component are comparable in the magnetized (bx50 and bx100) and unmagnetized (bx0) runs. The negligible role of the field in providing vertical support in the gas warmer than 200 K is consistent with the phase diagrams in Figure 14: the magnetic pressure is the highest fraction of the total pressure in the cold, dense gas.



**Figure 14.** Panels (a)–(e) represent pressure components ( $P/k$ ) plotted vs. density. Panel (a): total pressure (sum of thermal, kinetic, and magnetic). Panel (b): thermal pressure ( $nT$ ). The gray line shows the curve of thermal equilibrium ( $n^2\Lambda = n\Gamma$ ) for our cooling curve and diffuse heating function at  $z = 0$  for  $T < 2 \times 10^4$  K. Panel (c): turbulent pressure ( $(\rho)\sigma_{\text{turb}}^2/k$ , calculated in boxes 8 pc on a side for  $|z| < 125$  pc). The black dots show turbulent pressure calculated in boxes 125 pc on a side. Panel (d): kinetic pressure ( $(\rho)|v|^2/k$ ). Panel (e): magnetic pressure ( $(\mathbf{B})^2/8\pi k$ ). Panel (f) shows the ratio of thermal pressure to magnetic pressure (the plasma  $\beta$ ). All panels except (c) were calculated for every cell in the range  $|z| < 20$  pc of a snapshot of the bx50hr model at 340 Myr.

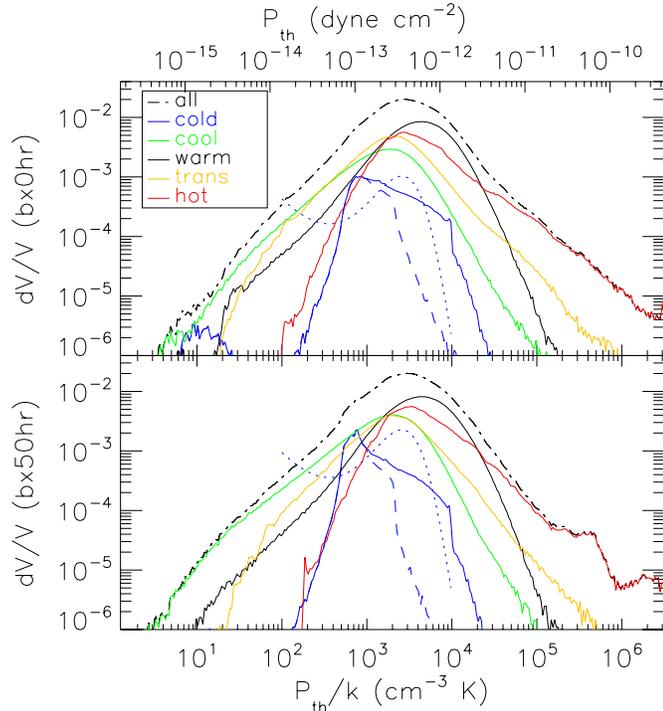
The scale heights of the two dominant (by mass) phases of the simulated ISM are closely related to the scale heights of the physical processes we imposed upon the simulation (summarized in Table 3). The cold gas scale height (8 pc in the unmagnetized model and 10–11 pc in the magnetized models) is comparable to the characteristic scale height of the gravitational acceleration near the plane ( $z \ll z_0$ ), where  $a_1 + a_2 = 8.77 \text{ pc}^{-1} (\text{pc Myr}^{-2})^{-1}$ . We have written the units here to emphasize that the  $a_1$  and  $a_2$  parameters in Equation (7) are the linear coefficients describing the change in gravitational acceleration (units of  $\text{pc Myr}^{-2}$ ) with height (units of pc). The warm gas scale height (190–200 pc in all models) is comparable to the core-collapse supernova (90 pc) and diffuse heating (300 pc) scale heights. Because the warm gas is thermally stable primarily due to the balance of radiative cooling and diffuse heating (the gray line in Figure 14(b)), the warm gas scale height is likely tied more strongly to the diffuse heating than to the supernovae.

In our models, the vertical support is provided primarily by turbulent and thermal pressure, with magnetic pressure playing only a small role in even the magnetized runs (Section 3.6 and Figures 16 and 17). This contrasts with models that include

rotational shear with (Gressel et al. 2008a) or without (Piontek & Ostriker 2007) supernovae; in these models, the magnetic pressure does provide significant vertical support. Without supernova driving turbulence, the role of turbulent pressure is small (Piontek & Ostriker 2007).

#### 4.4.1. Transition-temperature Gas

The near-unity volume- and mass-filling fractions of transition-temperature gas at  $|z| \approx 1$  kpc (dashed lines in Figures 10 and 11) deserve special attention. Because of the high cooling efficiency of O VI and other high ions, gas at  $T \sim 10^5$  K is normally assumed to be in transition, either cooling or heating, at an interface between hot and warm layers. The images of density and temperature at  $|z| = 1$ –2 kpc (Figure 3) appear strongly mixed, with slightly less mixing in the magnetized cases than in the unmagnetized case. We also note the distributions of temperature within the  $10^{4.5} - 10^{5.5}$  K regime at  $|z| \sim 500$  pc (green lines in Figure 5). In the unmagnetized model, the most probable temperature is  $\approx 10^{4.8}$  K, a local minimum in our cooling curve (Figure 1 of JM06), whereas the most probable temperature is  $\approx 10^{5.5}$  K in the magnetized case. This is consistent with previous numerical work that found large amounts of



**Figure 15.** Volume-weighted histograms of thermal pressure. The bx0hr (top panel) and bx50hr (bottom panel) models in the range  $|z| < 20$  pc are shown averaged over 21 time slices from 320 to 340 Myr. Gas in each temperature range is identified by color in the legend. The dashed blue line shows only cold gas from our simulation with  $n < 100 \text{ cm}^{-3}$ . The distribution identified in C I data for the CNM by Jenkins & Tripp (2011) is shown with the blue dotted line, with the maximum normalized to the maximum value of  $dV/V$  for the cold gas distribution in our simulation. (The dotted line is not otherwise fit to our results.) Pressures are labeled in both  $P/k$  units (bottom axis) and cgs units (top axis). Histograms were calculated with logarithmic bin intervals  $d(\log P/k) = 0.01$ .

O VI-absorbing gas in supernova cavities in the plane (Avillez & Breitschwerdt 2005b); they argued that turbulent mixing produces a large number of interfaces over a large volume. The suppression of turbulent mixing by the field thus reduces the efficiency of cooling in a turbulent plasma.

In this interpretation, the transition-temperature gas is shock-heated and cooling. Our simulations do not include Lagrangian tracer particles, so we cannot identify the past history of transition-temperature gas. However, because our diffuse heating is turned off at temperatures  $T > 2 \times 10^4$  K, the only heating mechanism capable of producing  $10^5$  K gas is supernovae and the shocks they produce, suggesting that the widespread transition-temperature gas at  $|z| = 1\text{--}2$  kpc is most likely heated by repeated shocks propagating upward from the plane.

The observed scale heights of ions that trace transition-temperature gas, determined by comparing the vertical component of the column densities of these ions to the height above the plane of the stars toward which the column densities are measured, are  $\sim 3$  kpc (Savage & Wakker 2009). This is more than three times the 0.6–1 kpc scale height in our models, found by directly measuring the vertical distribution of the transition-temperature gas. A direct comparison of synthetic and real observations will provide the best test of whether this apparent disagreement is real but is beyond the scope of this paper.

#### 4.4.2. Oscillation

In all cases, the scale heights oscillate. This behavior was predicted in one-dimensional models by Walters & Cox (2001).

In their models, compressions near the midplane drive oscillations in the vertical extent of the ISM. Material near the plane has a “bounce time” of 40–60 Myr, while the bounce time for higher material is  $\sim 100$  Myr. This behavior can be seen in the images of the evolving halo (Figure 2), in which gas within a few kiloparsecs of the plane rises initially and returns to the disk after 60–90 Myr, whereas the hot shocks which extend to  $\sim 10$  kpc return to the disk after  $\sim 100\text{--}200$  Myr. Comparison of the scale heights in Figure 12 shows similar behavior: an initial bounce of the hot gas (red lines) returns to a minimum (the most-compressed mode of Walters & Cox) after 120 Myr, while the warm gas (black lines) oscillates with a shorter period. We reiterate that our star formation rate is constant, so no additional star formation is triggered in the compressed mode. In addition, the amplitude of the oscillations is damped in the magnetized case. This is particularly obvious in the cold gas, which is also the most magnetically dominated.

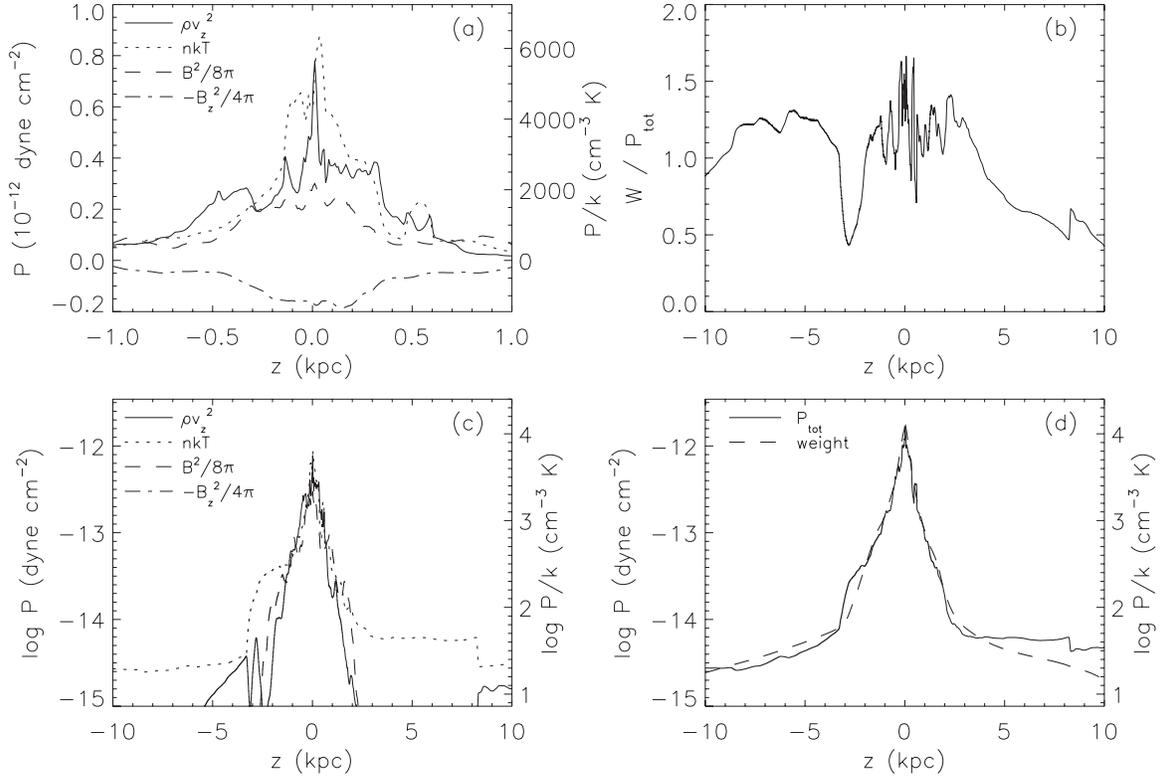
The Walters & Cox (2001) model also provides a framework for understanding the delayed second set of shocks we discussed in Section 3.1. After the end of the first oscillation of the gas close to the plane at  $t \approx 100$  Myr, a significant column of gas continues to fall from above, leading to a stationary shock which we see at  $|z| \approx 3$  kpc. Once a sufficient column density has accumulated near this shock, a second oscillation is triggered. This occurs at  $t \approx 150$  Myr in our models.

#### 4.5. Energy and Generation of the Magnetic Field

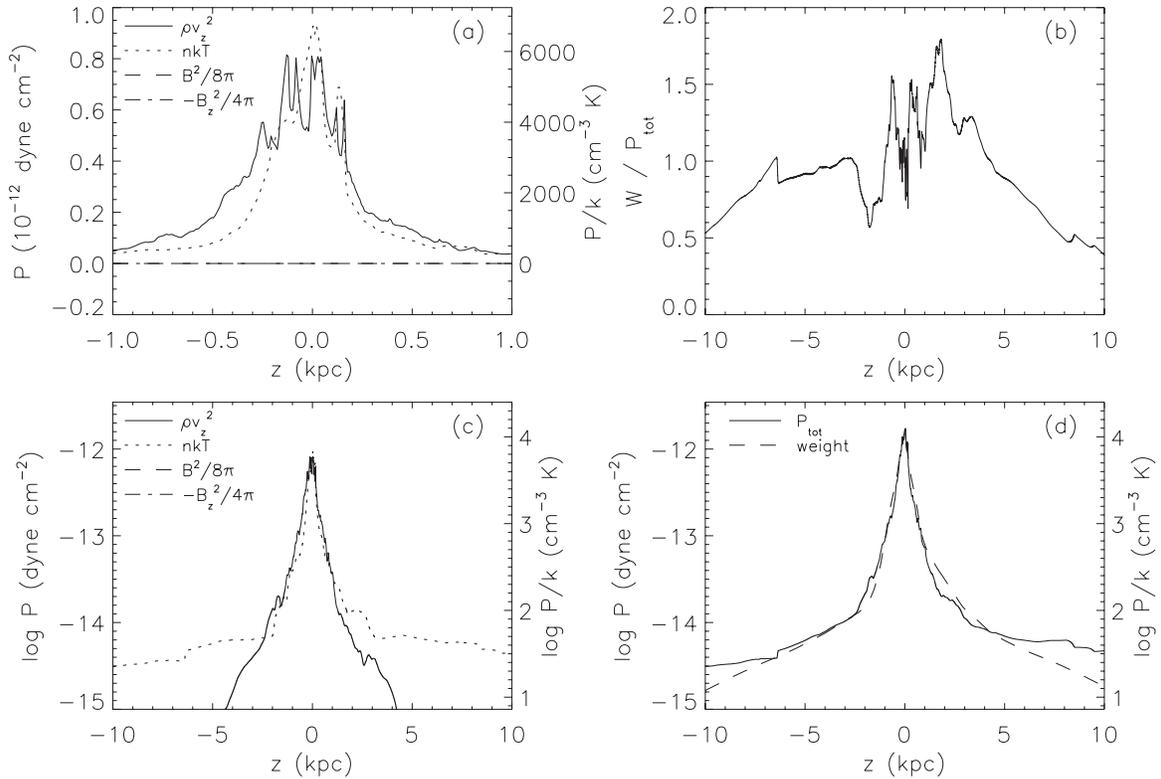
The total energy (kinetic, thermal, and magnetic) contained within our box is constant to within  $\sim 20\%$  after 20 Myr, suggesting that the turbulence is well developed. We now focus our discussion on the bx0hr and bx50hr models because the bx100 model is run at low resolution and its field had not reached a statistical steady state by the time the run finished after 260 Myr. The decrease in magnetic energy over time (Figure 13) is likely due to numerical dissipation, with a floor in the magnetic energy established by a turbulent dynamo. The kinetic and thermal energies in the bx0 and bx50 are comparable throughout the runs (as shown for the thermal energy in Figure 13). The total energy in the bx50 model is larger than that in the unmagnetized run, with the difference in total energy comparable to the initial magnetic energy.

Most measurements indicate that the magnetic field in the Milky Way consists of an ordered, azimuthal field of  $1.5\text{--}2.0 \mu\text{G}$  with an rms field of  $5\text{--}6 \mu\text{G}$  (see review by Kulsrud & Zweibel 2008). After 340 Myr, the mass-weighted rms magnetic field strength in the bx50hr model is  $4 \mu\text{G}$ , compared to the initial value of  $\mathbf{B}_{\text{ini}} = (6.5, 0, 0) \mu\text{G}$ . The field remains partially aligned with the initial field: the mass-weighted mean magnetic field vector after 340 Myr is  $\mathbf{B} = (2.0, -0.2, -0.2) \mu\text{G}$ , with standard deviations of  $(3, 4, 2) \mu\text{G}$ .

The physical mechanism responsible for maintaining the interstellar field in the Galaxy is still being debated. In supernova-driven simulations in a fully periodic box without gravity, Balsara et al. (2004) found that the magnetic field amplifies with an  $e$ -folding time of  $\sim 10$  Myr once turbulence is fully developed. Their simulations did not run for sufficient time to reach a steady state magnetic energy density. To test similar growth, we have run an additional model at 4 pc resolution, bx10, noting that Sur et al. (2010) and Federrath et al. (2011) have demonstrated effective dynamos with the HLL3R Riemann solver used in FLASH in this work. This model began with magnetic energy at 4% of the thermal energy, similar to the final state of the Balsara et al. (2004) work. The magnetic and thermal energy



**Figure 16.** Panel (a): kinetic ( $\rho v_z^2$ ), thermal ( $nkT$ ), magnetic ( $B^2/8\pi$ ), and magnetic tension ( $-B_z^2/4\pi$ ) pressures for the  $t = 340$  Myr snapshot of the  $bx50hr$  model at  $|z| < 1$  kpc. Values are space-averaged over  $x$  and  $y$  for each height  $z$ . Panel (b): ratio of total weight to total pressure. Panel (c): same as panel (a) but with a logarithmic vertical axis and expanded to show  $|z| < 10$  kpc; the magnetic tension is omitted because of its negative value. Panel (d): total pressure  $P_{tot}(z)$  (dashed line) and weight  $W(z)$  (solid line, as defined in Equation (11)) as a function of height.



**Figure 17.** Same as Figure 16 but for the  $t = 340$  Myr snapshot of the unmagnetized  $bx0hr$  model.

densities of the bx10 model are shown as a function of time in Figure 13 (dashed and solid red lines, respectively). Unlike the models with stronger initial fields, the field in the bx10 model grows initially (after turbulence is established). The field reaches a maximum surface magnetic energy density of  $\approx 0.3 \text{ erg pc}^{-2}$  after 40 Myr; the total magnetic energy is approximately 18% of the thermal energy in the box at this point. After 170 Myr, the magnetic energy falls, leaving the magnetic energy at 8% of the thermal energy. The reason for this behavior is unclear.

The decrease of the magnetic energy in the bx100 and bx50 models, which each have initial magnetic energies comparable to or larger than the thermal energy, combined with the increase (compared to the initial condition) of the magnetic energy in the bx10 case suggests that the field tends toward a preferred fraction of equipartition. However, the field energy after  $\sim 300$  Myr of evolution is dependent upon the initial field energy at 4 pc resolution. Noting that Balsara et al. (2004) used a resolution of  $\approx 1$  pc in their supernova-driven dynamo model, a high-resolution version of the bx10 model may prove illuminating.

Given the decrease in the magnetic energy to  $\lesssim 1/3$  of equipartition in our models, it is clear that a dynamo driven by supernovae alone will not create a magnetic field with an important role in the vertical stratification of the ISM. Rotational shear can amplify the field exponentially to at least  $0.5 \mu\text{G}$  (Figure 3 of Gressel et al. 2008a), though no interstellar dynamo model to date has been run for long enough to allow the field to reach equipartition. In their models without supernovae, Piontek & Ostriker (2007) found that a dynamo driven by differential rotation allows the magnetic energy density to reach 0.3–0.6 of equipartition with the thermal energy density, but the turbulent energy density is quite small absent supernovae.

#### 4.6. Numerical Resolution

Our approach of running the simulations at medium (4 pc) resolution for 300 Myr before switching to high (2 pc) resolution suggests that our main results are well converged. The mass fractions (Figure 7) and vertical distributions (Figure 12) in each component of the ISM are modified by  $< 1\%$  between the medium- and high-resolution runs, indicating that the overall distribution of material, particularly as a function of height, is well converged at the 4 pc midplane resolution. The volume-filling fractions in the plane (Figure 8) are affected more significantly, with the volume occupied by cold and warm gas each decreasing by  $\approx 10\%$  in the 5 Myr immediately following the resolution switch. Both volume-filling fractions then reach a steady state at the new resolution. Cold clouds exist at sizes down to the resolution limit at both medium and high resolutions, indicating that we do not resolve the smallest clouds. However, such collapse may not be physically meaningful without including self-gravity and modeling of the formation of molecules.

#### 4.7. Missing Large-scale Effects

We model an approximation of physical parameters appropriate for the solar circle. With these parameters, the simulations reach a statistical steady state only after  $\gtrsim 200$  Myr in the plane, with oscillations in the halo continuing for the entire run of our simulations. However, several factors related to Galactic structure on scales larger than 1 kpc are missing from our work.

*Spiral arms.* The supernova rate adopted here corresponds roughly to a star formation rate averaged over the Galactic

disk. In reality, massive stars form in spiral arms which pass a given parcel of the Galaxy every  $\sim 200$  Myr. Therefore, a more complete model would vary the core-collapse supernova rate.

*Accretion.* Our model does not include accretion. High-velocity clouds (HVCs), which have typical masses of  $\sim 10^6\text{--}10^7 M_\odot$ , provide a mass influx of  $\sim 1 M_\odot \text{ yr}^{-1}$  over the entire Galactic disk (Wakker et al. 2007; Thom et al. 2008). The interaction of HVCs with the Galactic halo likely has a significant effect on the behavior of the ISM and may drive turbulence. However, the mechanism by which infalling material joins the Galactic ISM is unknown and is beyond the scope of this paper.

*Differential rotation.* Differential rotation may be an important driver of turbulence, particularly in the outer Galaxy through the magnetorotational instability (Piontek & Ostriker 2007). Its inclusion may allow some form of large-scale dynamo to function, as discussed in Section 4.5.

## 5. CONCLUSIONS

We have considered the role of supernovae in the vertical stratification of components of a magnetized ISM in collisional ionization equilibrium in the Galactic gravitational field using stable, positivity-preserving 3D AMR MHD simulations. We summarize our main conclusions here.

1. Our models contain distinct cold and warm phases of the ISM (Figures 5, 6, and 14(b)). Reasonable treatment of the diffuse heating rate, which models photoelectric heating of gas cooler than  $2 \times 10^4$  K, is necessary for distinct phases to be present. With excessive diffuse heating, stable equilibria do not exist (JM06).
2. The qualitative structure of the ISM is similar in magnetized and unmagnetized models (Figure 3). The field in our models also has very little effect on the vertical distribution or scale heights of the components of the ISM (Figure 9). This is because the horizontally averaged magnetic pressure ( $|\mathbf{B}|^2/2$ ) and magnetic tension force ( $-B_z^2$ ) nearly cancel each other out at all heights (Figure 16(a)). Even ignoring magnetic tension, the magnetic pressure dominates over the thermal pressure only in the cold gas, and the cold gas is the only phase which shows some evidence for an increased scale height in the magnetized model. This effect is small because kinetic pressure is dominant in the cold gas in our models (Figures 14(d) and (a)).
3. Hydrostatic equilibrium is maintained primarily by thermal and kinetic pressure. We do not confirm the hypothesis of Wood et al. (2010) that magnetic fields would help to address the smaller-than-observed scale height of the WIM in the hydrodynamical models used in that work. In both cases, we do not model photoionization heating, which could lead to additional vertical support for the WIM. Rotational shear may well be important in allowing the magnetic field to provide vertical support (Piontek & Ostriker 2007; Gressel et al. 2008b).
4. The addition of a field reduces the mass in cold gas by 8% with a corresponding increase in the mass contained in lower density, warm gas (Figure 7), consistent with our expectation that the magnetic field would reduce the compressibility of the gas. The volume-filling fraction of the cold gas in the plane is also  $\approx 20\%$  higher in the 2 pc resolution, magnetized model than in the similar unmagnetized one (thick blue lines in Figure 8 and blue lines in Figure 11).

5. Extending the box to heights  $|z| \gtrsim 10$  kpc to fully capture the Galactic fountain flow is the single most important numerical factor in establishing a realistic temperature distribution in the halo (Section 2.1.2).
6. Even though our star formation rate does not change over time, the gas oscillates vertically (Figures 2 and 12), as predicted by Walters & Cox (2001).
7. At heights  $1 \text{ kpc} \lesssim |z| \lesssim 2 \text{ kpc}$ , the modeled ISM consists primarily of thermally unstable “transition-temperature” ( $T \sim 10^5$  K) gas (Figure 10). This is a thick transition region between the disk—mostly warm gas with cold clouds and hot supernova remnants—and the hot halo. In this transition region, the gas is strongly mixed, with interfaces between hot and warm gas filling most of the volume, leading to the near-unity filling fraction of  $10^5$  K gas.

The authors thank K. Wood for helpful discussions and E. Jenkins for a comment on the C I distribution. A.S.H. thanks S. Friedman and B. Morsony for sharing their expertise with FLASH, his thesis committee (M. A. Bershad, J. S. Gallagher, D. McCammon, and E. M. Wilcots) for comments on a draft of this paper, and the American Museum of Natural History for hospitality during the initial stages of this work. The FLASH code was developed by the DOE-supported ASC/Alliance Center for Astrophysical Thermonuclear Flashes at the University of Chicago. The computations were performed using National Science Foundation-supported TeraGrid resources provided by the Texas Advanced Computing Center under grant No. TG-MCA99S024. This work was partly supported by NASA/SAO grant TM0-11008X, NASA ATP grant NNX10A170G, and NSF grants AST-1109395 and AST-0607512.

## REFERENCES

- Avillez, M. A. 2000, *MNRAS*, 315, 479
- Avillez, M. A., & Breitschwerdt, D. 2004, *A&A*, 425, 899
- Avillez, M. A., & Breitschwerdt, D. 2005a, *A&A*, 436, 585
- Avillez, M. A., & Breitschwerdt, D. 2005b, *ApJ*, 634, L65
- Bakes, E. L. O., & Tielens, A. G. G. M. 1994, *ApJ*, 427, 822
- Balsara, D. S., Kim, J., Mac Low, M., & Mathews, G. J. 2004, *ApJ*, 617, 339
- Boulares, A., & Cox, D. P. 1990, *ApJ*, 365, 544
- Bregman, J. N. 1980, *ApJ*, 236, 577
- Brunt, C. M., Heyer, M. H., & Mac Low, M.-M. 2009, *A&A*, 504, 883
- Chiang, W.-H., & Prendergast, K. H. 1985, *ApJ*, 297, 507
- Clemens, D. P., Sanders, D. B., & Scoville, N. Z. 1988, *ApJ*, 327, 139
- Cox, D. P., & Smith, B. W. 1974, *ApJ*, 189, L105
- Dalgarno, A., & McCray, R. A. 1972, *ARA&A*, 10, 375
- Dehnen, W., & Binney, J. 1998, *MNRAS*, 294, 429
- Dickey, J. M., & Lockman, F. J. 1990, *ARA&A*, 28, 215
- Federrath, C., Sur, S., Schleicher, D. R. G., Banerjee, R., & Klessen, R. S. 2011, *ApJ*, 731, 62
- Fellhauer, M., Belokurov, V., Evans, N. W., et al. 2006, *ApJ*, 651, 167
- Ferrière, K. M. 2001, *Rev. Mod. Phys.*, 73, 1031
- Flynn, C., Holmberg, J., Portinari, L., Fuchs, B., & Jahreiß, H. 2006, *MNRAS*, 372, 1149
- Fox, A. J., Wakker, B. P., Savage, B. D., et al. 2005, *ApJ*, 630, 332
- Fryxell, B., Olson, K., Ricker, P., et al. 2000, *ApJS*, 131, 273
- Gómez, G. C., Benjamin, R. A., & Cox, D. P. 2001, *AJ*, 122, 908
- Gressel, O., Elstner, D., Ziegler, U., & Rüdiger, G. 2008a, *A&A*, 486, L35
- Gressel, O., Ziegler, U., Elstner, D., & Rüdiger, G. 2008b, *Astron. Nachr.*, 329, 619
- Habing, H. J. 1968, *Bull. Astron. Inst. Neth.*, 19, 421
- Haffner, L. M., Reynolds, R. J., & Tufté, S. L. 1999, *ApJ*, 523, 223
- Heiles, C., & Troland, T. H. 2003, *ApJ*, 586, 1067
- Heiles, C., & Troland, T. H. 2005, *ApJ*, 624, 773
- Helmi, A. 2004, *ApJ*, 610, L97
- Henley, D. B., Shelton, R. L., Kwak, K., Joung, M. R., & Mac Low, M.-M. 2010, *ApJ*, 723, 935
- Hill, A. S., Benjamin, R. A., Kowal, G., et al. 2008, *ApJ*, 686, 363
- Hollenbach, D. J., Werner, M. W., & Salpeter, E. E. 1971, *ApJ*, 163, 165
- Jenkins, E. B., & Tripp, T. M. 2011, *ApJ*, 734, 65
- Johnston, K. V., Law, D. R., & Majewski, S. R. 2005, *ApJ*, 619, 800
- Joung, M. K. R., & Mac Low, M.-M. 2006, *ApJ*, 653, 1266 (JM06)
- Joung, M. R., Bryan, G. L., & Putman, M. E. 2012, *ApJ*, 745, 148
- Joung, M. R., Mac Low, M.-M., & Bryan, G. L. 2009, *ApJ*, 704, 137 (JMB09)
- Kainulainen, J., Beuther, H., Banerjee, R., Federrath, C., & Henning, T. 2011, *A&A*, 530, A64
- Kainulainen, J., Beuther, H., Henning, T., & Plume, R. 2009, *A&A*, 508, L35
- Kennicutt, R. C., Jr. 1989, *ApJ*, 344, 685
- Klingenberg, C., Schmidt, W., & Waagan, K. 2007, *J. Comput. Phys.*, 227, 12
- Klingenberg, C., & Waagan, K. 2010, *Acta Math. Sci.*, 30, 621
- Korpi, M. J., Brandenburg, A., Shukurov, A., Tuominen, I., & Nordlund, Å. 1999, *ApJ*, 514, L99
- Koyama, H., & Ostriker, E. C. 2009a, *ApJ*, 693, 1316
- Koyama, H., & Ostriker, E. C. 2009b, *ApJ*, 693, 1346
- Kuijken, K., & Gilmore, G. 1989, *MNRAS*, 239, 605
- Kulsrud, R. M., & Zweibel, E. G. 2008, *Rep. Prog. Phys.*, 71, 046901
- Law, D. R., & Majewski, S. R. 2010, *ApJ*, 714, 229
- Mac Low, M.-M., & Ferrara, A. 1999, *ApJ*, 513, 142
- Mac Low, M.-M., & Klessen, R. S. 2004, *Rev. Mod. Phys.*, 76, 125
- Mac Low, M.-M., McCray, R., & Norman, M. L. 1989, *ApJ*, 337, 141
- Madsen, G. J., Reynolds, R. J., & Haffner, L. M. 2006, *ApJ*, 652, 401
- Manchester, R. N., Fan, G., Lyne, A. G., & Kaspi, V. M. 2006, *ApJ*, 649, 235
- Matzner, C. D. 2002, *ApJ*, 566, 302
- McKee, C. F. 1990, in *ASP Conf. Ser. 12, The Evolution of the Interstellar Medium*, ed. L. Blitz (San Francisco, CA: ASP), 3
- McKee, C. F., & Ostriker, J. P. 1977, *ApJ*, 218, 148
- Navarro, J. F., Frenk, C. S., & White, S. D. M. 1996, *ApJ*, 462, 563
- Norman, C. A., & Ferrara, A. 1996, *ApJ*, 467, 280
- Parravano, A., Hollenbach, D. J., & McKee, C. F. 2003, *ApJ*, 584, 797
- Piontek, R. A., & Ostriker, E. C. 2007, *ApJ*, 663, 183
- Reynolds, R. J., & Cox, D. P. 1992, *ApJ*, 400, L33
- Reynolds, R. J., Haffner, L. M., & Tufté, S. L. 1999, *ApJ*, 525, L21
- Rosen, A., & Bregman, J. N. 1995, *ApJ*, 440, 634
- Rosen, A., Bregman, J. N., & Kelson, D. D. 1996, *ApJ*, 470, 839
- Savage, B. D., & Wakker, B. P. 2009, *ApJ*, 702, 1472
- Shapiro, P. R., & Field, G. B. 1976, *ApJ*, 205, 762
- Sur, S., Schleicher, D. R. G., Banerjee, R., Federrath, C., & Klessen, R. S. 2010, *ApJ*, 721, L134
- Sutherland, R. S., & Dopita, M. A. 1993, *ApJS*, 88, 253
- Thom, C., Peek, J. E. G., Putman, M. E., et al. 2008, *ApJ*, 684, 364
- Vázquez-Semadeni, E., & García, N. 2001, *ApJ*, 557, 727
- Waagan, K. 2009, *J. Comput. Phys.*, 228, 8609
- Waagan, K., Federrath, C., & Klingenberg, C. 2011, *J. Comput. Phys.*, 230, 3331
- Wakker, B. P., York, D. G., Howk, J. C., et al. 2007, *ApJ*, 670, L113
- Walters, M. A., & Cox, D. P. 2001, *ApJ*, 549, 353
- Wolfire, M. G., Hollenbach, D., McKee, C. F., Tielens, A. G. G. M., & Bakes, E. L. O. 1995, *ApJ*, 443, 152
- Wolfire, M. G., McKee, C. F., Hollenbach, D., & Tielens, A. G. G. M. 2003, *ApJ*, 587, 278
- Wood, K., Hill, A. S., Joung, M. R., et al. 2010, *ApJ*, 721, 1397