# An Introduction to Generalized Probabilistic Theories 

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## Generalized Probabilistic Theories

Generalized Probabilistic Theories

- Framework to describe measurement statistics of general physical theories
- Generalization of the density matrix formalism in quantum theory

Motivation
■ Why quantum theory? (search for new physical theories)

- Alternatives to quantum theory


## Generalized Probabilistic Theories

Usual experimental setup

state $\omega$
measurement M with outcomes i

## Generalized Probabilistic Theories

Usual experimental setup


## Definition (State space)

The state space $\Omega$ is the set of all states $\omega$ allowed in the theory

## Generalized Probabilistic Theories

Usual experimental setup


## Definition (Effects)

Effects are maps $e_{i}^{M}: \Omega \rightarrow[0,1]$ that give probabilities of measurement outcome $i$ when measuring state $\omega$ :

$$
p(i \mid \omega, M)=e_{i}^{M}(\omega)
$$

The set of all effects is $E$

## Minimal physical requirements

■ Probabilistic applications of preparation and measurement devices:

1) Convexity:

$$
\begin{aligned}
& \forall \omega=\sum_{i} \lambda_{i} \omega_{i}, \lambda_{i} \geq 0, \sum_{i} \lambda_{i}=1, \omega_{i} \in \Omega: \omega \in \Omega \\
& \forall e=\sum_{i} e_{i} \omega_{i}, \lambda_{i} \geq 0, \sum_{i} \lambda_{i}=1, e_{i} \in E: e \in E
\end{aligned}
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$\Rightarrow \Omega, E$ are convex sets

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## Definition (Pure states)

Pure states are given by extremal points in $\Omega$

## Minimal physical requirements

- Probabilistic applications of preparation and measurement devices:

2) Linearity (choice of devices should be independent of their measurment statistics):

$$
\begin{aligned}
e\left(p \omega_{1}+(1-p) \omega_{2}\right) & =p e\left(\omega_{1}\right)+(1-p) e\left(\omega_{2}\right) \\
{\left[p e_{1}+(1-p) e_{2}\right](\omega) } & =p e_{1}(\omega)+(1-p) e_{2}(\omega)
\end{aligned}
$$

$\Rightarrow e, \omega$ are elements of linear spaces $A, A^{*}$

## Minimal physical requirements

■ Operationalism: Equivalence of objects leading to the same measurement statistics

$$
\begin{aligned}
& e\left(\omega_{1}\right)=e\left(\omega_{2}\right) \forall e \in E \Rightarrow \omega_{1}=\omega_{2} \\
& e_{1}(\omega)=e_{2}(\omega) \forall \omega \in \Omega \Rightarrow e_{1}=e_{2}
\end{aligned}
$$

## Minimal physical requirements

■ Completeness: all linear mappings $e: \Omega \rightarrow[0,1]$ included as valid measurement outcomes in a theory
$\Rightarrow \exists$ unique effect $u: u(\omega)=1 \quad \forall \omega \in \Omega$ called the unit measure

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## Definition (Measurement)

A measurement $M$ is a set of effects $\left\{e_{i}^{M}\right\}$ summing up to the unit measure $u=\sum_{i} e_{i}^{M}$ :

$$
\Rightarrow \sum_{i} p(i \mid \omega, M)=\sum_{i} e_{i}^{M}(\omega)=\left[\sum_{i} e_{i}^{M}\right](\omega)=u(\omega)=1
$$

## Positive Cones

- State space: vectors on a hyperplane $u(\omega)=1$

■ Unnormalized states form cone $A_{+}$
$\epsilon$


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- State space: vectors on a hyperplane $u(\omega)=1$

■ Unnormalized states form cone $A_{+}$
■ Effects: Elements of dual cone $A_{+}^{*}$


## Examples

## Classical Probability Theory

■ $\Omega$ is a simplex spanned by $\operatorname{dim}^{A}$ linear independent extremal points $\omega_{i}$
■ For each pure state $\omega_{i}$ exists an extremal effect $e_{i}$ that identifies it uniquely: $e_{i}\left(\omega_{j}\right)=\delta_{i j}$
$\omega_{i}=\left(\begin{array}{c}\cos \left(\frac{2 \pi(i-1)}{}\right) \\ \sin \left(\frac{2 \pi(i-1)}{4}\right) \\ 1\end{array}\right)$
$e_{i}=\omega_{i}$
$u=(1,1,1)^{T}$

$\mathrm{n}=3$

■ $\omega_{i}$ form basis for $A \Rightarrow$ unique decomposition of mixed states only for classical systems

## Examples

Quantum theory
■ $A_{+}=A_{+}^{*}$ is the cone of positive hermitian matrices $\rho \geq 0, \rho^{\dagger}=\rho$

- $e(\rho)=\operatorname{tr}[e . \rho] \quad u=\mathbb{1}$
$\Rightarrow \Omega=\left\{\rho \in A_{+} \mid u(\rho)=\operatorname{tr}[\rho]=1\right\}$
Qubit: $\Omega$ is given by 3-dim Bloch sphere


## Examples

The Gbit
■ $\Omega$ given by a square


■ No uncertainty for pure states with respect to measurements $M_{1}=\left\{e_{1}, e_{3}\right\}$ and $M_{2}=\left\{e_{2}, e_{4}\right\}$

## Joint Systems

Definition (Global State Assumption)
Joint system $A B$ fully characterized by joint probabilities $\left\{p\left(e^{A}, e^{B}\right)\right\}$

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Global State Assumption
No-Signalling Principle

$$
\} \Longrightarrow A B_{+} \quad \subseteq \quad A \otimes B
$$

## Cone Of Joint Systems

■ Cone $A B_{+}$of joint system bounded:

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A_{+} \otimes_{\min } B_{+} \subseteq A B_{+} \subseteq A_{+} \otimes_{\max } B_{+}
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A_{+} \otimes_{\min } B_{+}=\text {ConvexSpan }\left\{\omega^{A} \otimes \omega^{B}\right\}
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A_{+} \otimes_{\max } B_{+}=\left\{\omega^{A B} \in A \otimes B \mid \omega^{A B}\left(e^{A} \otimes e^{B}\right) \geq 0\right\}
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includes all joint states allowed by No-Signalling
- Theorie defined by structure of $A_{+}, u^{A}, B_{+}, u^{B}$ and $A B_{+} \cdots$


## Duality relations

For finite dimensional systems:

$$
\begin{aligned}
\left(A_{+} \otimes_{\max } B_{+}\right) & =\left(A_{+}^{*} \otimes_{\min } B_{+}^{*}\right)^{*} \\
\left(A_{+} \otimes_{\min } B_{+}\right) & =\left(A_{+}^{*} \otimes_{\max } B_{+}^{*}\right)^{*}
\end{aligned}
$$

■ All tensor products equal or there exists at least either states or effects that are entangled
■ In QT: Balanced between entangled effects and states

## Thank you for your attention!

